

## VAST: An ASKAP Survey for Variables and Slow Transients

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**Abstract:** The Australian Square Kilometre Array Pathfinder (ASKAP) will give us an unprecedented opportunity to investigate the transient sky at radio wavelengths. In this paper we present VAST, an ASKAP survey for Variables and Slow Transients. VAST will exploit the wide-field survey capabilities of ASKAP to enable the discovery and investigation of variable and transient phenomena from the local to the cosmological, including flare stars, intermittent pulsars, X-ray binaries, magnetars, extreme scattering events, interstellar scintillation, radio supernovae and orphan afterglows of gamma ray bursts. In addition, it will allow us to probe unexplored regions of parameter space where new classes of transient sources may be detected. In this paper we review the known radio transient and variable populations and the current results from blind radio surveys. We outline a comprehensive program based on a multi-tiered survey strategy to characterise the radio transient sky through detection and monitoring of transient and variable sources on the ASKAP imaging timescales of five seconds and greater. We also present an analysis of the expected source populations that we will be able to detect with VAST.

**Keywords:** telescopes — surveys — radio continuum: general — galaxies: general — stars: general — ISM: general

# 1 Introduction

Many classes of astronomical objects are known to be variable radio sources, including the Sun, the planets, cool stars, stellar binary systems, neutron stars, supernovae, gamma-ray bursts and active galactic nuclei. There have been extensive studies of these objects using targeted surveys, but few blind, unbiased surveys for variable radio phenomena. An effective survey for radio variability needs to maximise the metric

$$A\Omega \left( \frac{T}{\Delta t} \right) \quad (1)$$

where  $A$  is the effective collecting area of the telescope,  $\Omega$  is the solid angle coverage of the survey,  $T$  is the total duration of the observations and  $\Delta t$  is the time resolution (Cordes et al. 2004; Cordes 2007). Typically blind radio surveys have had either too small a field of view and too little time coverage per field of view to allow a comprehensive survey for transient and variable phenomena.

The Square Kilometre Array (SKA)<sup>1</sup> is a proposed future radio telescope (Dewdney et al. 2009) that will be many times more sensitive than any existing facility. Should it go ahead, the SKA will be designed to explore five key science projects: galaxy evolution, cosmology and dark energy; strong-field tests of gravity using pulsars and black holes; the origin and evolution of cosmic magnetism; probing the dark ages — the first black holes and stars; and the cradle of life, searching for life and planets. In addition, there is an awareness of building in (or not designing out) the possibility of serendipitous discoveries. In particular, the dynamic radio sky is recognised as a rich and relatively unexplored discovery space (Cordes et al. 2004; Cordes 2007). Plans for the SKA are still under development, with a Phase 1 instrument expected to be completed in 2020, and a Phase 2 instrument in 2024.

The Australian Square Kilometre Array Pathfinder (ASKAP)<sup>2</sup> is a precursor and technology development platform for the full SKA and will greatly improve on our ability to detect radio transients and variable sources. Its moderately high sensitivity (RMS of  $\sim 1$  mJy beam<sup>-1</sup> in 10 seconds), combined with a wide field of view (covering 30 square degrees of sky in a single pointing) will enable fast, sensitive all-sky surveys (Johnston et al. 2007, 2008). ASKAP is currently under construction in Western Australia. The final telescope will have 36 12-metre antennas, with a maximum baseline of 6 km. The wide field of view is enabled by the use of Phased Array Feeds (Chippendale et al. 2010) of 100 dual-polarisation pixels that will operate over the range 700–1800 MHz. A summary of the planned ASKAP specifications is given in Table 1. The 6-antenna Boolardy Engineering Test Array (BETA), is expected to be completed in late 2012 and will allow some early testing of the techniques required for the full ASKAP surveys.

ASKAP is primarily a survey instrument, and this is reflected in the time allocation policy. In its initial

Table 1: ASKAP Specifications

Number of dishes	36
Dish diameter (m)	12
Dish area (sq m)	113
Total collecting area (sq m)	4072
Aperture efficiency	0.8
System temperature (K)	50
Field-of-view (deg <sup>2</sup> )	30
Frequency range (MHz)	700–1800
Bandwidth (MHz)	300
Maximum number of channels	16 384
Maximum baseline (km)	6
RMS cont. sensitivity (10 s)	640 $\mu$ Jy/bm
RMS cont. sensitivity (1 hr)	47 $\mu$ Jy/bm

five years of operation 75% of the available time will be scheduled for several major Survey Science Projects<sup>3</sup> that were selected through a competitive call for proposals (Ball 2009).

In this paper we describe one of these projects, An ASKAP Survey for Variables and Slow Transients (VAST)<sup>4</sup>. In the context of ASKAP, ‘slow’ means variability or transient behaviour detectable on timescales greater than the cadence on which images will be performed (expected to be 5–10 seconds). A related project, The Commensal Real-Time ASKAP Fast-Transients Survey (CRAFT; Macquart et al. 2010), will search for ‘fast’ transients on timescales shorter than 5 seconds. Different techniques are required to process and analyse data for fast and slow transients, and from a scientific perspective this division roughly represents a separation between coherent and incoherent emission processes. The main science goals of VAST are:

- to detect and monitor ‘orphan’ gamma-ray burst afterglows to understand their nature;
- to conduct an unbiased survey of radio supernovae in the local Universe;
- to determine the origin and nature of the structures responsible for extreme scattering events;
- to make a direct detection of baryons in the intergalactic medium;
- to discover flaring magnetars, intermittent or deeply nulling radio pulsars, and rotating radio transients through changes in their pulse-averaged emission;
- to investigate intrinsic variability of active galactic nuclei;
- to detect and monitor flare stars, cataclysmic variables and X-ray binaries in our Galaxy; and
- to discover previously unknown classes of objects.

To achieve these goals VAST will conduct three surveys, in addition to commensal observing with other

<sup>1</sup>See <http://www.skatelescope.org>.

<sup>2</sup>See <http://www.atnf.csiro.au/projects/askap>.

<sup>3</sup>The full list of projects is available at <http://www.atnf.csiro.au/projects/askap/ssps.html>.

<sup>4</sup>See <http://www.askap.org/vast>.

Table 2: Survey parameters for the VAST surveys (see Section 4.1 for details). The two components of the VAST-Deep survey are discussed further in the text. The Commensal column gives an indication of the amount of time that VAST will have access to through commensal observing with other ASKAP projects.

	VAST-Wide	VAST-Deep		VAST-Galactic	Commensal
		Multi-field	Single field		
Observing time (hrs)	4380	3200	400	600	1.5 years
Survey area (sq deg)	10 000	10 000	30	750	10 000
Time per field	40 s	1 hr	1 hr	16 min	12 hours
Repeat	daily	7 times	daily	64 times	none
Observing freq (MHz)			1130–1430		
Bandwidth (MHz)			300		
RMS sensitivity ( $\text{mJy beam}^{-1}$ )	0.5	0.05		0.1	0.01
Field of view (sq deg)			30		
Angular resolution			10''		
Spectral resolution			10 MHz		
Time resolution			5 seconds		
Polarisation products			IQUV		

ASKAP projects: VAST-Wide, aimed at detecting rare, bright events, will survey 10 000 square degrees per day to an RMS sensitivity of  $0.5 \text{ mJy beam}^{-1}$ ; VAST-Deep, aimed at detecting unknown source classes, will survey 10 000 square degrees down to an RMS sensitivity of  $50 \mu\text{Jy beam}^{-1}$ ; and VAST-Galactic which will cover 750 square degrees of the Galactic plane, plus the Large and Small Magellanic Clouds down to an RMS sensitivity of  $0.1 \text{ mJy beam}^{-1}$ . The full survey parameters are shown in Table 2 and discussed in Section 4.

In Section 2 we motivate the VAST surveys by discussing the science goals in more detail, and explain how they fit in the context of existing efforts. We also make predictions about the likely source populations that will be observed by VAST. In Section 3 we review previous and current blind radio surveys, as well as archival studies. In Section 4 we describe the proposed VAST survey strategy in detail. Finally, in Section 5 we discuss the design of the VAST transient detection pipeline.

## 2 VAST Science

Astronomical transient phenomena are diverse in nature, but they can be broadly classified on the basis of their underlying physical mechanism into one of four general categories: explosions, propagation effects, accretion-driven, and magnetic field-driven events. In this section we discuss the main VAST science goals in each of these areas and discuss the expected populations for different types of sources.

### 2.1 Explosions

Gamma-ray bursts (GRBs) and supernovae (SNe) are some of the most energetic explosions in the Universe. Although they have been studied extensively there are

a number of unresolved questions about their nature, including the existence of radio orphan afterglows (Levinson et al. 2002), the beaming fraction of gamma-ray bursts (Sari et al. 1999; Dalal et al. 2002), and the current rate of massive star formation in the Universe as traced by the cosmological radio supernova rate. These questions can be addressed by surveys that observe a large area of sky to good sensitivity, on a regular basis. The capacity of VAST to detect these rare objects is orders of magnitude greater than previous blind surveys such as those of Gal-Yam et al. (2006) or Bower et al. (2007).

#### 2.1.1 Gamma-Ray Bursts

The afterglow emission from gamma-ray bursts is generated via particle acceleration in the decelerating blast wave. At late times, the peak of the afterglow spectrum is at radio energies, and thus radio studies provide the most interesting probes for GRBs science. Based on previous studies (Chandra & Frail 2012), we expect that the GRBs seen by ASKAP will be long-lasting but intrinsically fainter on average than shorter wavelength afterglows. Given the sensitivity of ASKAP, this means that those GRBs studied by VAST will, in general, be the rare but important nearby ( $\sim 1500 \text{ Mpc}$ ) events. In the relativistic expansion phase (up to 100 days after the initial explosion), ASKAP monitoring in conjunction with optical and X-ray facilities will allow us to measure physical parameters such as the total kinetic energy of the explosion, the density of the circumburst medium and the geometry of the outflow.

At late times, when the blast wave has decelerated to sub-relativistic velocities, the jetted outflows expand and become more isotropic. Under these circumstances, the the total (i.e. calorimetric) energy of GRB blast waves may be estimated - a key but difficult parameter to measure (Frail et al. 2005). Instead of relying on rare, nearby events, VAST may also obtain a robust beaming-independent estimates GRB en-

ergies from a large sample of single-epoch flux density measurements (and limits) of afterglows at late times (Shivvers & Berger 2011).

As the relativistic beaming is totally suppressed at late time, radio monitoring is the best solution to identify ‘orphan afterglows’, i.e., those GRBs whose jet is beamed away from us (Rhoads 1997). While they would lack a gamma-ray trigger by definition, such systems could be distinguished from other classes of transients through the evolution and luminosity of their light curves, their location towards galaxies that resemble other GRB hosts, and their probable association with supernovae (Soderberg et al. 2010). The discovery of a population of orphan afterglows would allow us to constrain the beaming fraction of GRBs – a key quantity for understanding the true GRB event rate. It would also open up an order of magnitude more objects which can be scrutinised in detail. This may also give insights on the jet structure of GRBs and clarify the puzzling lack of jet breaks from the *Swift* satellite (Racusin et al. 2009).

### 2.1.2 Supernovae

While thermonuclear (Type Ia) supernovae have proven to be extremely useful in measuring the acceleration of the Universe, no Type Ia SN has ever been detected at radio wavelengths (Panagia et al. 2006; Hancock et al. 2011; Chomiuk et al. 2012), indicating a very low density of circumstellar material. In contrast, many core-collapse SNe (those of Type Ib, Ic, and the various Type II sub-classes) have been detected and monitored at radio wavelengths (Weiler et al. 2002; Ryder et al. 2004; Soderberg 2007). Radio observations replay the late phases of stellar evolution (in reverse), and reveal clues about the late-time evolution of the progenitor system.

With VAST we will conduct an unbiased census of core-collapse SNe (a still undetermined fraction of which emit at radio wavelengths), allowing us to match radio detections against optically-discovered SNe. We will also detect new radio SNe that may have gone undetected in the optical or infrared due to significant amounts of dust (Kankare et al. 2008). The prevalence of such heavily-obscured SNe is virtually unconstrained, with hints that they may constitute a non-trivial subpopulation. Recently, SN 2008cs (Kankare et al. 2008) was detected in the IR and confirmed in the radio bands, and SN 2008iz (Brunthaler et al. 2009) was serendipitously discovered in radio observations of M 82. SN 2008iz has remained undetected in the optical, IR, and X-ray bands despite sensitive searches (Brunthaler et al. 2010). Discovery of new SNe is valuable for tying down the current rate of massive star formation in the Universe and for understanding how and when SNe explode in star-bursting galaxies. Radio emission from core-collapse SNe is detectable within days, and they reach their peak luminosity at 1.4 GHz as much as one year after the explosion, making them ideal targets to detect and study with the sensitivity and cadence of VAST (see Section 4).

Beyond the large population of standard core-collapse SNe, there exist sub-classes of objects that are of par-

ticular interest. For instance, radio observations of a recent type Ib/c supernova, SN2009bb, showed that it was expanding at mildly relativistic velocities ( $\beta \approx 0.8$ ) powered by the central engine (Soderberg et al. 2010). Such SNe are usually associated with GRBs and are discovered by their intense but short-lived gamma-ray energy emission. Yet, SN2009bb had no detected gamma-ray counterpart. It has been suggested that the blast-waves produced in such engine-driven SNe could be responsible for the acceleration of ultra-high energy cosmic-rays (Chakraborti et al. 2011). However, the present limits on the neutrino flux expected as a by-product of cosmic-ray production in GRBs are at least a factor of 3.7 below predictions (Abbasi et al. 2012). Radio surveys with the capability of completing the sample of these explosions will improve our understanding of particle acceleration processes and the diversity of core-collapse events (see Lien et al. 2011, for predictions on the detection rates of core-collapse radio SNe for a range of SKA pathfinder instruments).

### 2.1.3 Gravitational Wave Counterparts

Gravitational wave (GW) events are expected to be compact and involve significant energies—both GRBs and SNe are considered to be among likely sources of both electromagnetic and GW emission. In an analogous manner, other classes of GW events may also generate significant transient radio emission, especially if they occur in dense environments. We discuss the particular case of mergers of neutron star binaries in Section 2.3.4. For all classes of GW events, however, localisation will be a significant challenge, with search areas expected to be of order 10–100 deg<sup>2</sup> for planned GW detectors (Fairhurst 2011). The large instantaneous sky coverage of ASKAP makes it particularly valuable in the search for electromagnetic counterparts to GW events, although there are significant uncertainties in the requisite search sensitivities and timescales. The search area estimates mentioned above are computed for the Advanced LIGO (Harry et al. 2010) and Advanced VIRGO (Accadia et al. 2011) detectors. These systems are planned to come online in the 2014–2015 timeframe, meaning that ASKAP will be well-positioned to make discoveries in this parameter space.

## 2.2 Propagation

**Plasma Propagation Effects:** At centimetre wavelengths any object with an angular size  $\lesssim 1$  mas is subject to the effects of interstellar scintillation (ISS), caused by inhomogeneities in the ionised interstellar medium (ISM) of our Galaxy. Refractive ISS (RISS) is manifested as slow intensity variations with a wide correlation bandwidth that is caused by large-scale refracting structures in the ISM. Much faster variations that have narrow frequency scales are associated with diffraction from small-scale irregularities (DISS). Observations of the intensity fluctuations caused by ISS can be used to infer both the small-scale structure of the ISM and the structure of a scintillating source on scales down to  $\sim 10 - 200$  microarcseconds for compact extragalactic sources and  $\sim 10$  nanoarcseconds



for some pulsars.

The density inhomogeneities in the ionised ISM are typically characterised as a power law (Armstrong et al. 1995). However, a small number of pulsars have displayed strong fringing events in their dynamic spectra, in which the normal, random appearance of the dynamic spectrum caused by DISS is modified by the appearance of quasi-periodic fringes. Both pulsars and AGN have displayed extreme scattering events (ESEs) in their light curves, in which their flux densities exhibit marked decreases ( $\approx 50\%$ ) bracketed by modest increases; and a number of AGN have exhibited intraday variability (IDV), marked by flux density variations on time scales of hours. Fringing events are interpreted most naturally as the beating of two distinct images of the pulsar and both ESEs and IDV can be interpreted as due to lensing events through a plasma. All of these phenomena could occur if there is a temporary increase in the amount of power on scales of order  $10^{13}$  cm along or close to the line of sight, e.g., as due to a discrete “cloud” of electrons drifting in front of the pulsar or AGN (Cordes & Wolszczan 1986; Clegg et al. 1998). In the case of IDV, it may also be the case that transient, high brightness temperature structures are required to be produced in the cores of some AGN.

The characteristics of the discrete structures inferred to be responsible for these phenomena are such (density  $> 100 \text{ cm}^{-3}$ , temperature  $> 10^3 \text{ K}$ ) that they would be highly overpressured with respect to the bulk properties of the ISM (Kulkarni & Heiles 1988), and would likely be transient structures. While there is some evidence to suggest that a small fraction of the volume of the ISM is at a high pressure (Jenkins & Tripp 2001), the distribution of such highly overpressured structures is not known, although the characteristics of IDV in some AGN suggest that some structures are extremely close to the Sun ( $\sim 1 - 10 \text{ pc}$ ). The existence of such highly over-pressured structures in the local ISM is hard to reconcile with our current understanding of interstellar turbulence (e.g. Fiedler et al. 1987b; Walker 2007). Alternatively, these phenomena may indicate that a simple power law is not a good description of the spectrum of density inhomogeneities, which would then constrain the nature of the plasma processes involved in the generation and maintenance of the density inhomogeneities.

Because these propagation events are transient in nature, many of the previous observations have been “accidental,” happening to catch a source displaying such an event during the course of observations designed for other reasons. There have been relatively few surveys designed specifically to try to find propagation events, with notable examples being the Green Bank Interferometer (GBI) monitoring program (Fielder et al. 1994; Lazio et al. 2001b) and the MASIV survey (Lovell et al. 2003, 2008). Moreover, because of either limited sensitivity, limited field of view, or both, most surveys have only been able to monitor a relatively small number ( $\sim 100$ ) of relatively strong sources ( $\gtrsim 100 \text{ mJy}$ ).

A wide-area, high cadence, sensitive survey would be able to place strong constraints on both the spatial distribution of the structures responsible for prop-

agation events and the likelihood of a source exhibiting such an event. In turn, these much stronger constraints can be used to infer properties of the ionised ISM and the mechanisms by which such structures are generated. As an illustration of the kind of monitoring program that ASKAP will enable, a survey could be designed to produce daily flux density measurements of sources stronger than about  $10 \text{ mJy}$  with the capability to detect  $50\%$  flux density changes at the  $10\sigma$  level. At this sensitivity level, a survey of even only  $1000 \text{ deg}^2$  would still be able to monitor nearly compact  $50,000$  sources.

**Gravitational Lensing:** Aside from the plasma propagation effects discussed in more detail below, we also note that strong gravitational lensing of active galaxies by galaxy clusters can lead to an achromatic time-domain signal that contains information about both the active galaxy and about the intervening spacetime. Radio observations of strong lensing have been conducted successfully with existing facilities and those data have been used to constrain cosmological parameters (e.g. Fassnacht et al. 1999; Myers et al. 2003). These operate by spatially separating the lensed components of the active galaxy. In general ASKAP will not have the angular resolution to observe such multiple imaging, but the large number of sources monitored with regular cadence means that auto-correlation analysis (as performed by Barnacka et al. 2011) can itself constrain lens time delays.

### 2.2.1 Pulse Broadening from Scattering

The well known temporal broadening of pulsar pulses caused by scattering from electron-density variations will affect any source of short duration bursts, including those that will be targeted by VAST. Pulse broadening from sources in the Galactic center can be broader than  $1000 \text{ sec}$  at  $1 \text{ GHz}$ . Bursts from extragalactic sources at cosmological distances can similarly be broadened, especially if the line of sight intersects another galaxy but also from the intergalactic medium (IGM). Scattering in the Milky Way, in other galaxies, and in the IGM therefore define pulse-broadening horizons beyond which pulses with specific widths will not be seen (see, e.g., Cordes 2012). Pulse broadening from scattering can be distinguished from source-intrinsic effects by identifying across the ASKAP band the strong frequency dependence  $\sim \nu^{-4}$  of scatter broadening.

### 2.2.2 Extreme Scattering Events

Extreme scattering events (ESEs) are flux density variations marked by decreases in flux density bracketed by increases (Figure 1). There is widespread agreement that ESEs are caused by refractive lens-like density perturbations along the line-of-sight to a background source (e.g. Fiedler et al. 1987b; Romani et al. 1987). The refracting lenses which cause ESEs must be Galactic in origin and probably lie within a few kpc, but in the 24 years since their discovery, no satisfactory physical model has emerged for the ESE phenomenon. Analysis of existing data suggests that ESEs are unlike any other identified component of the ISM and

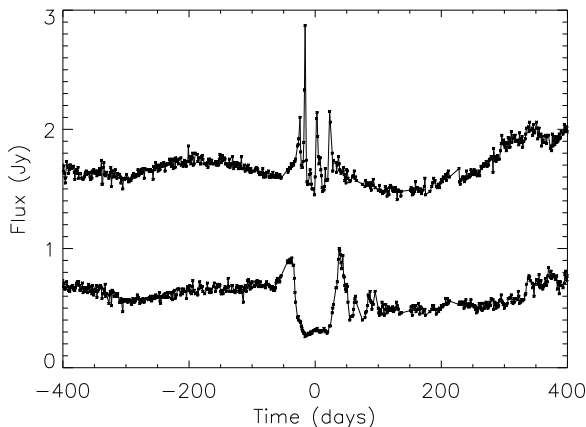


Figure 1: An Extreme Scattering Event in Q0954+658 at 2.7 GHz (lower) and 8.1 GHz (upper) adapted from Fiedler et al. (1987b). For clarity, an offset of 1 Jy has been added to the top trace. The strong frequency dependence of ESEs and the necessity of regular sampling of the light curves over a long period is evident. At ASKAP frequencies ( $\sim 1$  GHz), the amplitude of the flux density decrease would likely have been even larger, and the flux density would likely have increased to an even higher value during the start of the event.

that they are present in vast numbers in the Galaxy (with a density  $\sim 10^3 \text{ pc}^{-3}$ ). If the refracting properties are taken at face value, they imply regions with pressures  $\sim 1000$  times higher than the diffuse ionized ISM. One suggestion is that they are associated with neutral, self-gravitating gas clouds (Walker 2007), in which case they would make a substantial contribution to the total baryonic mass content of the Galaxy. Alternatively, they may represent small-scale density enhancements created by converging flows in a turbulent medium (Jenkins & Tripp 2001), but it is unclear whether such structures could persist for the observed ESE durations.

A wide-area survey, with daily cadence would allow us to address the physical composition and origin of the objects responsible for ESEs in a comprehensive manner.

**Population statistics:** One of the more comprehensive ESE surveys (Fiedler et al. 1994b) provided a total coverage of approximately 600 source-years; VAST has the potential to achieve a coverage that is larger by a factor of roughly 300, searching for ESEs toward sources with much lower flux densities. In turn, increased coverage will constrain the spatial density and distribution of refracting lenses in the Galaxy, including whether they are associated with particular Galactic structures (Fiedler et al. 1994a). A larger number of ESEs will also allow investigation into any potential diversity in light curve shapes, and whether any such diversity can be understood in terms of simple models (Clegg et al. 1998) or the extent to which more complex models of the density structures, background sources, or both are required.

**Real-time characterisation of ESEs:** A straightforward approach to elucidating the nature of the structures responsible for ESEs would be to conduct additional observations of a source while it is undergoing an ESE. Events like those seen towards quasars B0954+658 and B1749+096 (Fiedler et al. 1994b) or B1741–038 (Fiedler et al. 1992; Clegg et al. 1996) should be seen at the rate of roughly 60 per year in an all-sky monitoring survey with  $\sim 1$  mJy RMS sensitivity, with  $\sim 10$  in progress at any given moment. Thus, in contrast to the GBI monitoring program, for which only one ESE was identified while it was on-going, VAST will enable real-time follow-up.

Some key follow-up measurements could be made at higher frequencies, particularly in the optical to X-ray bands, which should directly reveal the dense neutral gas clouds responsible for ESEs via scattering, absorption and refraction (Draine 1998). Additionally, H I observations would reveal any associated neutral structures, Faraday rotation measurements may constrain the magnetic field strength within the structures, and high resolution observations could confirm predicted changes in the apparent structure of the source. While there have been previous efforts in conducting such observations (Clegg et al. 1996; Lazio et al. 2000, 2001a), only one source has been observed in such a manner (B1741–038), and, for some of the measurements, only a single observation during the actual ESE could be obtained. An all-sky survey could detect multiple ESEs during the course of a year, allowing ample opportunity to trigger follow-up observations with other telescopes and at other wavelengths.

### 2.2.3 Interstellar Scintillation of AGN

Observations of decimetre variability in extragalactic sources (e.g., Hunstead 1972) led quickly to the recognition that the sources needed to have extremely compact components, be affected by interstellar scintillation, or both (Armstrong et al. 1977; Spangler et al. 1989). The need for extremely compact components, apparently violating the so-called inverse Compton catastrophe, became even more acute with the discovery of “flickering” (Heeschen 1984), and later intra-day variability.

The MASIV 5GHz VLA Survey (Lovell et al. 2003, 2008) and other observations (Kedziora-Chudczer et al. 1997; Dennett-Thorpe & de Bruyn 2002; Bignall et al. 2003, 2006) have established that the fast—*intra-day* and *inter-day*—flux density variations at centimetre wavelengths exhibited by a subset of compact quasars are predominantly caused by ISS rather than intrinsic variability. *Intra-day* variability refers to sources whose characteristic timescale is less than 24 hours, while *inter-day* variability refers to sources whose timescale exceeds this.

Scintillating AGN are of astrophysical interest because the small angular sizes that they must possess in order to exhibit rapid ISS require brightness temperatures near, or possibly in excess of, the  $10^{12}$  K inverse Compton limit for incoherent synchrotron radiation. In several cases over the past decade, it ap-

peared that the Doppler boosting factors required in order to reconcile the apparent brightness temperature with the inverse Compton limit were well in excess of the typical values determined from Very Long Baseline interferometric observations of quasars. Estimates of the Doppler boosting factors are subject to considerable uncertainty related to the distance of the scattering material from Earth, but many of the most rapid IDVs appear to require Doppler factors of a few tens (Macquart et al. 2000; Rickett et al. 2002; Bignall et al. 2006; Macquart & de Bruyn 2007).

More recently, it appears that the fastest manifestations of IDV are associated with sources whose lines of sight intersect more local, albeit much more inhomogeneous, patches of turbulence than previously supposed. Thus there has been a transfer of our ignorance away from the physics of the AGN and onto the physics of interstellar turbulence. However, this issue is not yet conclusively resolved, as little is known about the origin (or physical plausibility) of the local, highly inhomogeneous scattering screens that need to be invoked to explain the observed scintillation properties of these sources.

The four-epoch MASIV survey (Lovell et al. 2003, 2008) detected variability in 56% of the 482 sources observed but the characteristics of the variability changed from epoch to epoch. While able to illustrate this ISS intermittency, the MASIV survey itself, because of its limited duration and cadence, could not address many of the resulting questions, notably the extent to which the intermittency results from changes in the ISM properties or changes in the sources. In order to address the cause of the intermittency, a survey must have the following characteristics. It must observe a large number of sources widely distributed on the sky so as to have robust statistics and to be able to search for correlations expected for ISS (e.g., a Galactic latitude dependence or with nearby structures in the ISM). The survey must measure both the total and linear polarization intensity of the sources—ISM-driven intermittency should result in the linear polarization scintillation tracking that of the total intensity, while source-driven intermittency should result in the linear polarization and total intensity scintillating nearly independently of each other. The survey must be conducted with close to daily cadence to capture the scintillation variations, and over a multi-year time scale in order to separate changes in source structure or scattering medium from changes associated with the annual cycle resulting from the Earth’s orbital velocity.

In addition to determining the cause of ISS intermittency and mapping the scattering properties of the local ISM, such a survey would address the existence of AGN emission above the inverse Compton brightness limit and potentially constrain the baryonic content of the IGM as a function of redshift.

**The cause of ISS intermittency:** By continuously observing a large sample of sources over a long period, VAST will comprehensively determine the intermittency properties of scintillating AGN. It will discriminate between ISM-related and source-driven effects via polarization variations, as discussed above. A

complication arises because annual cycles in the variability timescale (caused by the Earth’s velocity vector changing with respect to the scattering screens) need to be distinguished from true IDV intermittency. However, these two phenomena can be separated by making regular (daily) observations over the course of at least two years. The variations should be dissimilar between epochs separated by a year if the scintillation is intermittent. Intermittency, if caused by the ISM, points to the physical nature of the turbulent patches responsible for IDV.

**AGN emission above the inverse Compton brightness limit:** In some cases fast scintillation may be attributable to the appearance of micro-arcsecond structure within a source, associated with an outburst (Macquart & de Bruyn 2007). The rise and fade times of the manifestation of fast variations allow us to determine the longevity of micro-arcsecond components associated with IDV: The rate at which the AGN supplies energy to power the bright emission that causes the variability, and the rate at which energy losses cause its eventual decay. These rates can then be compared with specific AGN processes to identify the mechanism associated with the super-Compton radio emission.

**Baryonic content of the intergalactic medium:** The MASIV survey showed that few sources above  $z \sim 2$  are variable (Lovell et al. 2008). If this apparent suppression of ISS as a function of source redshift is confirmed by VAST, there are a number of possible explanations: (1) intrinsic source evolution; (2) gravitational microlensing; or (3) scattering in the turbulent ionised IGM. If the last is at least partially responsible, then the angular broadening is caused by the cumulative effect of all the baryons in the ionised IGM. The sheer number of sources VAST can monitor will enable us to probe in detail the evolution of structure in the ionised IGM as a function of redshift.

One potential complication with scintillation studies is that the largest amplitude IDV is typically observed at frequencies near 5 GHz. At ASKAP frequencies, it is possible that the strength of scattering, even toward high Galactic latitudes, and intrinsic source diameters will result in large-amplitude scintillations being quenched (Dennison & Condon 1981), although lower amplitude refractive scintillations may still be exhibited (e.g., Dennison et al. 1987).

## 2.3 Accretion and Magnetism

Along with explosions and propagation effects, accretion and magnetic fields are frequently implicated in transient events. Accretion onto a compact object is a powerful energy source that drives both local sources such as X-ray binaries (XRBs) and microquasars as well as extragalactic blazars and AGN at cosmological distances. The accretion flow is often connected to an outflow in the shape of a jet, which is the source of the radio emission. The intimate link between the accretion disc and the jet is demonstrated by the fundamental plane of black hole activity that describes a correlation between the mass of the black hole, the jet

power, and the accretion power (Merloni et al. 2003; Falcke et al. 2004) in active galactic nuclei. This relation can be extended to lighter black holes where it matches the correlation between the X-ray and radio luminosities in XRB black hole candidates (BHCs) that was found by Gallo et al. (2003). Radio emission is also observed from XRBs containing weakly magnetised neutron stars, such as Z and atoll sources, although they are at least one order of magnitude fainter in radio than the BHCs (e.g. Fender 2006), and a weak radio jet has also been observed from the cataclysmic variable SS Cyg during a dwarf nova outburst (Körding et al. 2008). There is still not a good theoretical understanding of how accretion discs produce outflows, but most models assume that the jet is accelerated by a magnetic field (e.g. Blandford & Payne 1982). VAST will allow us to characterise the full range of transient jet phenomena across the fundamental plane.

Magnetic fields can also act in other ways. They can convert the rotational energy of their source into electromagnetic energy and the mechanical energy of an outflow, as happens for instance in a radio pulsar, or they can release their own energy in the form of flares on a star or the radio emission from Jupiter and (possibly) Jupiter-like extrasolar planets. With VAST we will be able to reveal the unifying physical principles that underlie such diverse behaviour.

### 2.3.1 Intrinsic Variability in Active Galactic Nuclei

As discussed in Section 2.2.3, Fiedler et al. (1987a) observed a sample of 33 AGN, with flux densities of a few Janskys, every day for more than 6 years. All sources were variable over this time span with typically smooth variations in flux with quasi-periods of months to years. VAST will produce light curves and detect this sort of variability for more than  $10^4$  sources over the whole sky. The *Fermi* gamma-ray satellite observes the whole sky every 3 hours and has detected blazar variability on timescales as short as 6 hours (Abdo et al. 2010). Simultaneous measurements of the gamma-ray and radio variability in blazars will provide us with a powerful tool for understanding the physics of the central engines in AGN.

Beyond single black holes in AGN, a number of binary supermassive black holes (SMBHs) are known or suspected (e.g., Komossa 2006; Rodriguez et al. 2006; Smith et al. 2010; Burke-Spolaor 2010). The evidence for these varies, but it is clear that with the mergers of galaxies the central SMBHs could also merge. The act of inspiralling may lead to radio transients accessible to ASKAP. This is because the SMBHs will likely move through a static magnetic field provided by an accretion disk, and this should drive electromagnetic jets (Palenzuela et al. 2010; Lyutikov 2011; Moesta et al. 2012). The detailed emission processes are unclear, but it is plausible that there could be a radio signature accessible to ASKAP before (O’Shaughnessy et al. 2011) or during (Kaplan et al. 2011) the merger. Numerous other electromagnetic mechanisms have also been proposed; see Schnittman (2011) for a detailed discus-

sion. Combining the radio observations with observations at other wavelengths or even with gravitational waves would significantly improve our understanding of galaxy mergers (e.g., Sesana et al. 2011).

### 2.3.2 Tidal Disruption Events

Tidal disruption events (TDEs) occur when a star wanders too close to a black hole at the centre of a galaxy or star cluster. The black hole tidal forces tear the star apart, but the debris continue on ballistic orbits around the star and are accreted as they approach periastron on the following orbit. The resulting sudden increase in the black hole accretion rate may be directly detected at optical and X-ray frequencies. After the initial flare up the luminosity decays as  $t^{-5/3}$  (Rees 1990). Several candidates of such TDEs have been detected serendipitously through all-sky X-ray and ultraviolet surveys (Grupe et al. 1999; Komossa & Greiner 1999; Komossa & Bade 1999; Esquej et al. 2007, 2008; Gezari et al. 2006, 2008, 2009; Lin et al. 2011; Gezari et al. 2012; Cenko et al. 2012a; Saxton et al. 2012).

Crucially, black hole accretion is often accompanied by relativistic jets, with a mechanical power often comparable to the radiative luminosity. The jet can be a source of beamed radio emission and hard X-rays in a way similar to a blazar. To a hard X-ray telescope such as the Burst Alert Telescope (BAT) on the *Swift* it will initially resemble a GRB, but it will be distinguished by its much longer duration. *Swift* J1644+57 is the first observed example of such a TDE candidate, at redshift  $z = 0.35$ . It was discovered as a luminous X-ray transient and was observed across the electromagnetic spectrum, enabling multi-band analysis and modeling of the event (Bloom et al. 2011; Levan et al. 2011; Zauderer et al. 2011; Berger et al. 2012; Wiersema et al. 2012). The mass of the central black hole responsible for this event was derived to be  $\lesssim 10^7 M_\odot$  and the ejecta Lorentz factor  $\gtrsim 10$ . More recently another TDE candidate has also been detected by Swift in hard X-rays, *Swift* J2058.4+0516 (Cenko et al. 2012b).

The main difference between AGN and TDE jets is that the former propagate to Mpc distances, and create characteristic cocoon/lobe structures where their kinetic power is dissipated. By contrast, TDE jets are impulsive ejections that produce a forward shock/reverse shock structure as they interact with the ISM (Sari & Piran 1995). As there is no pre-existing cavity around the previously quiescent black hole, TDE jets have to propagate through a much denser environment, and decelerate to sub-relativistic speed at distances  $\lesssim 1$  pc. Most of their kinetic power is expected to be dissipated in the reverse shock (Giannios & Metzger 2011). Typical TDE radio flares from the disruption of a solar-mass star by a  $10^7$  solar mass black hole are expected to peak in brightness after  $\sim 1$  year, reaching a flux density of  $\sim 2$  mJy at 1.4 GHz, for a distance of 1 Gpc corresponding to  $z = 0.2$  (not including the effect of relativistic boosting for jets oriented along our line of sight). At peak brightness, the expected radio spectrum is inverted (self-absorbed) with  $F_\nu \sim \nu^{-2}$  below  $\sim 2$  GHz, and  $F_\nu \sim \nu^{-1/3}$  above that



(Giannios & Metzger 2011).

Radio afterglows of TDEs would be a powerful tool to understand jet physics and source geometry, through measurement of the flux evolution and polarization. The expected degree of polarization for these events is low if they are collimated, similar to GRB afterglows. Recent infra-red and radio polarization measurements have been found to be consistent with theoretical expectations (Wiersema et al. 2012; Metzger et al. 2012).

Radio detections of TDE jets hold the key to other unsolved astrophysical problems (e.g. Krolik & Piran 2012). It is well known that not all accreting black holes have jets, but it is not clear why. Jets are formed only at certain ranges of accretion rate, but other parameters such as the black hole spin may also be important. Finding what fraction and what type of X-ray detected TDEs have radio jets will help address this question, and will shed light on the evolution of nuclear black holes in galaxies. Furthermore, the timescale for the formation of a radio-bright TDE jet will help us understand how quickly a collimated magnetic field is generated inside the transient accretion disk.

The predicted brightness of TDEs at gigahertz frequencies suggests that VAST should be able to detect about a dozen events per year. A systematic survey of the radio sky would be able to constrain the rate of TDEs and reduce the collimation uncertainty significantly. In particular, the discovery of an entirely new population of orphan TDEs would allow us to place solid constraints on the occurrence rate of this phenomenon and therefore test the theoretical prediction that we should see  $\sim 10^{-4}$  yr $^{-1}$  per galaxy (Rees 1990).

### 2.3.3 Accreting Neutron Stars, Black Holes and Microquasars

Exploring the connection between accretion disks and jets in X-ray binaries (XRBs) requires correlated radio and X-ray observations (Fender 2006). The VAST survey will reveal and subsequently measure the radio emission from a large number of XRBs and allow us to search for new radio outbursts. This, coupled with the recent availability of large field of view instruments like *Swift*-BAT and *Fermi*-GBM, and with the early results (from 2013) of eROSITA's all sky survey, will provide a large sample of XRBs with active jets, and will allow us to do statistical studies of their time variability properties and X-ray/radio correlation.

VAST will yield information on the population of XRBs showing quasi-steady state radio jets, which are mildly relativistic, and are associated with non-thermal, radiatively-inefficient accretion states (Fender et al. 2004; McClintock & Remillard 2006). For these persistent jets, fundamental properties such as baryonic content, energy budget, emission mechanism, and the role played by the spin of the compact object are not well constrained (Markoff et al. 2005; Narayan 1996; Yuan et al. 2005), and regular radio monitoring will allow us to investigate them. In addition, we will determine the duty cycle of the radio-loud phases, and test whether radio state transitions (changes in the kinetic power) are always associated with X-ray transitions (changes

in the radiative power). We will also investigate what physical mechanism causes major radio flares and the ejection of highly relativistic, optically-thin bullets, often associated to the state transition from the hard/non-thermal to soft/thermal state.

In addition to stellar mass and supermassive black holes, it is predicted that VAST will detect jets from intermediate mass black holes (IMBHs) with masses between  $\sim 100 - 100,000 M_{\odot}$ . IMBHs may provide a pathway for the formation of supermassive black holes (e.g. Ebisuzaki et al. 2001) and also have important connotations for other areas of astrophysics (e.g. Fornasa & Bertone 2008; Wang et al. 2010), but until recently their existence was highly disputed. The discovery of the object known as HLX-1 located in the galaxy ESO 243-49, with an X-ray luminosity  $\sim 1000$  times greater than the Eddington limit for a  $10 M_{\odot}$  black hole, currently provides very strong evidence for the existence of IMBHs (Farrell et al. 2009), with independent studies constraining the mass to between  $9,000 M_{\odot} < M_{BH} < 90,000 M_{\odot}$  (Davis et al. 2011; Servillat et al. 2011; Godet et al. 2012; Webb et al. 2012). Recently, transient radio emission was detected from HLX-1 following a transition from the low/hard to high/soft spectral states, consistent with the ejection of relativistic jets from an IMBH (Webb et al. 2012).

Objects such as HLX1 may be very rare (only one known within 100 Mpc), but there are hundreds of ultraluminous X-ray sources (ULXs; see Feng & Soria 2011 for a recent review) in the local Universe with luminosities exceeding the classical Eddington limit of stellar-mass black holes. Most of them are likely to be explained through other mechanisms such as anisotropic emission and/or super-Eddington accretion, but we cannot rule out that a few of them may be IMBHs. In addition, there may be hundreds of low-state or quiescent IMBHs lurking in the core of globular clusters, in the nuclei of dwarf galaxies, or inside recently merged galaxies that may occasionally become active. The detection of jet emission from these objects would allow us to infer the black hole mass, as the radio and X-ray luminosities are correlated as a function of mass (e.g. Fender 2006). Steady jets would be undetectable for all but the closest of these objects. However, transient radio flares have been observed from XRBs associated with jet ejection events with radio luminosities a factor of 10 – 100 times brighter than the non-flaring radio emission (e.g. Körding et al. 2006). Such flares should be detectable by VAST out to distances of  $\sim 5 - 15$  Mpc from  $1,000 M_{\odot}$  black holes, and out to  $\sim 23 - 72$  Mpc from  $10,000 M_{\odot}$  black holes (Farrell, in preparation). Thus we will be able to survey the local Universe for flaring IMBHs, placing constraints on the number density and therefore testing their validity as a pathway for supermassive black hole formation. Conversely, if ULXs are super-Eddington accretors, detection of radio emission will prove that even at those extreme accretion rates, a substantial fraction of accretion power can come out as kinetic power, and will improve our understanding of black hole feedback at high accretion rates.

### 2.3.4 Neutron Star Mergers

The merger of two compact objects (neutron stars) is expected to generate strong gravitational waves (§2.1.3) but they will also reveal themselves through electromagnetic radiation. The relativistic, beamed outflows detected from short-hard GRBs are thought to originate from such mergers (e.g Fong et al. 2012). VAST is not expected to fortuitously observe such beamed events. However, merger simulations also show that substantial mass can be ejected from these systems quasi-isotropically and at sub- or trans-relativistic velocities, and with energies of  $10^{49} - 10^{50}$  erg (Piran et al. 2012).

Consider a mildly relativistic blast wave ( $\beta_0 \sim 0.8$ ) moving out with an energy  $\sim 10^{49}$  erg into the surrounding medium. The shock-wave sweeps the surrounding medium and heats it to relativistic temperatures by converting bulk kinetic energy of the incoming material into thermal energy of the shocked material. The electrons in the shock-heated plasma are accelerated in the post-shock magnetic field to radiate synchrotron radiation. This would eventually be optically thin in the gigahertz regime, peaking on a timescale of about a couple of months when the blast wave starts to decelerate. The optimum waveband for detecting these transients, which are also supposed to be prime sources of gravitational waves, is around gigahertz frequencies (Nakar & Piran 2011). A merger within 300 Mpc would shine as a radio transient of about 1 mJy brightness at  $\sim 1$  GHz (Kamble & Kaplan in preparation) which would be detectable to the VAST and, potentially, also to Advanced-LIGO. Both the peak flux density and the time to peak are uncertain since we lack accurate estimates of the ejecta velocity and energy, and we expect the circumburst density to vary from  $1 \text{ cm}^{-3}$  to  $10^{-4} \text{ cm}^{-3}$ .

Using the ‘best-bet’ neutron star merger detection rates predicted for Advanced LIGO (Abadie et al. 2010), we estimate that up to about half a dozen of these transients should be detectable by ASKAP at any given time, although there are significant uncertainties involved. Piran et al. (2012) independently estimate ASKAP’s sensitivity to these events and obtain similar results. These are comparable to the rate estimated from observations of short GRBs which have been proposed to originate in similar mergers (Coward et al. 2012). Metzger & Berger (2012) emphasise the challenges related to uncertain GW event localisations, which are significant. As discussed in §2.1.3, the large instantaneous sky coverage of ASKAP may be a significant advantage in this situation.

### 2.3.5 Flaring and Pulsing Neutron Stars

Of the known neutron stars, the vast majority are pulsars that have been detected through their pulsed radio emission. VAST will not be sensitive to radio pulsations on timescales of  $\lesssim 1$  second, but will instead probe neutron star populations that are not selected in pulsar surveys. We will be sensitive to intermittent or deeply nulling radio pulsars as well as radio pulsars with large variations in their average pulsed flux.

Magnetars will be detectable through their rare flaring events such as the giant flare and expanding radio nebula produced by SGR 1806–20 (Gaensler et al. 2005; Cameron et al. 2005). In addition, some magnetars are known to turn on as radio sources. For example, XTE J1810–197 was initially detected as a transient radio source (Halpern et al. 2005) before its radio pulsations were identified (Camilo et al. 2006). We may also be able to detect bright intermittent pulses from the enigmatic rotating radio transients (RRATs; McLaughlin et al. 2006). VAST will thus reveal a broader census of the Galactic neutron star population than a standard pulsation search, with subsequent implications for the overall supernova event rate in our Galaxy.

The radio source FIRST J1023+0038 presents an interesting case study. Initially it was classified as a Galactic cataclysmic variable (CV) (Bond et al. 2002) because of its optical spectrum and variability, but its radio flaring behavior appeared to be anomalous. The classification was called into question because of the changes in the spectrum over time (e.g., Thorstensen & Armstrong 2005) and it was eventually re-identified as a transitional object where a recycled radio pulsar had turned on after a low mass X-ray binary phase (Archibald et al. 2009, 2010). The actual transition event had an extremely short duration, with the accretion disk apparently dissipating between 2001 December and 2002 May. The radio pulsar was probably absent in 1998 (upper limits of 1.8 and 3.4 mJy at 1.4 GHz) but detected in a survey in mid-2007 with a flux density  $\sim 14$  mJy at 1.6 GHz.

### 2.3.6 Classical novae

Novae occur in binaries (CVs) that consist of a white dwarf that is accreting matter from its companion star. A nova is a thermonuclear explosion in the layer of accreted matter on the surface of a white dwarf. A shell of matter is ejected during the explosion and this shell might produce radio emission. The radio emission can persist for a few years after the explosion, and in the case of FH Ser it reached  $\sim 10$  mJy at 2.7 GHz about one year after the outburst. In the late stages when the ejecta is optically thin, radio emission is generated by the entire ejecta and can therefore be used to estimate the total ejecta mass (see e.g., Hjellming 1990 for a review).

### 2.3.7 Flare stars, pre-main sequence stars, and rotation-driven stellar activity

There are several types of stars that display large flares that can be observed in the radio band. The most common are cool dwarf stars of spectral type M or later including brown dwarfs. The flares have typical flux densities above  $\sim 1$  mJy, but there is also quiescent emission at the sub-mJy level between the flares (e.g. Pestalozzi et al. 2000). For these types of systems it has been argued that radio observations are the best way to obtain information on the stellar magnetic field strength (Berger 2006). The number of flaring events that are potentially detectable is unclear, but Berger

(2006) obtained a detection rate of  $\sim 10\%$  for objects of later spectral class than M7. They also find that the radio to X-ray luminosity ratio increases over at least three orders of magnitude for these objects compared to earlier spectral classes. VAST will undertake a census of such stars in the local (10 pc) neighbourhood.

In AM Her-like systems (cataclysmic variables in which the white dwarf has a surface magnetic field of  $\sim 10^3 - 10^4$  T) the M dwarfs are forced to co-rotate with the binary, which has an orbital period of a few hours. The dwarf star then becomes more active; flares in AM Her itself have been observed to be as bright as 10 mJy at 5 GHz (Dulk et al. 1983).

Another group of stars that can produce radio flares are the RS CVn-systems (Gunn et al. 1994). These are close detached binaries consisting of two main sequence or evolved stars of spectral type F or later, that are tidally locked. The rapid spin of the stars leads to strong magnetic activity, which is manifested as large spots on the stars and also as flaring activity that can be observed both in X-rays and at radio wavelengths.

Flare-like outbursts are also observed in pre-main sequence stars, but it is not clear whether the physical mechanism is really that of a stellar flare or related instead to the accretion flow onto the star. Such outbursts have been observed in both naked T Tauri stars, which do not show signs of extended accretion discs, and in systems with an infrared excess, which indicates the presence of a disc (Osten & Wolk 2009).

A magnetic field can convert the rotational energy of a star into an outflow and electromagnetic radiation. The most well-known examples of this are the classical radio pulsars, but a similar phenomenon can arise in less extreme stars. Radio observations have detected two fully circularly polarised radio pulses per spin period from the Ap star CU Vir, which is also spinning down in a way similar to a pulsar (e.g. Trigilio et al. 2008; Trigilio et al. 2011; Lo et al. 2012).

Another example of activity driven by rotation is the white dwarf in the cataclysmic variable AE Aqr. The spin period of the white dwarf is 33 s, so rapid that the white dwarf does not accrete the matter that streams over from the companion K dwarf. Instead, the matter is ejected from the system. AE Aqr can be observed over a wide range of frequencies from X-rays to radio, and pulsed emission is observed at optical and X-rays, but it has not yet been detected at radio wavelengths (De Jager et al. 1994; Bookbinder & Lamb 1987; Bastian et al. 1988).

### 3 Exploration of the Unknown

The most interesting sources detected by VAST may be those we currently know nothing about. A sensitive wide-area blind survey is ideal for detecting such unknown source classes, but one of the challenges in designing the VAST survey is ensuring that the survey parameters do not unduly bias us against unknown source classes that we would otherwise be sensitive to. For example, daily monitoring of a field for 5 minutes per day makes it less likely that we will detect an object that brightens on hour-long timescales and fades

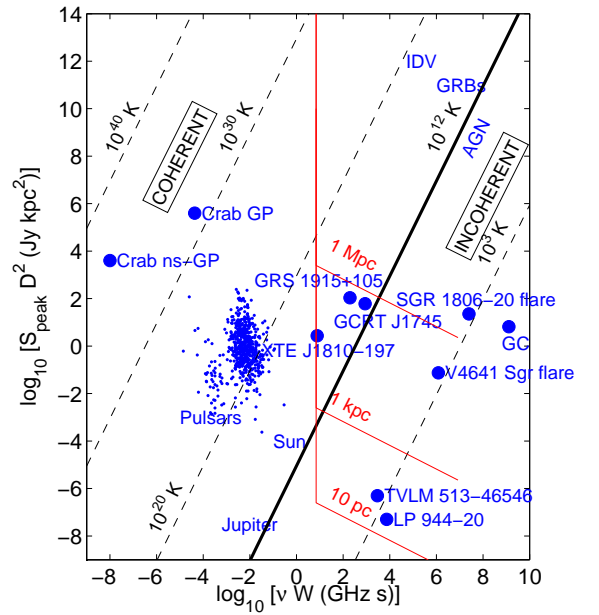


Figure 2: The parameter space for radio transients, adapted from Cordes et al. (2004). A quantity equivalent to absolute luminosity (observed flux density  $S$  multiplied by the square of the distance  $D^2$ ) is plotted against the dimensionless product of the emission frequency  $\nu$  and the transient duration or pulse width  $W$ . In the Rayleigh-Jeans approximation, these quantities are directly proportional and related to the brightness temperature, as indicated by the diagonal lines, with  $T = 10^{12}$  K marking the maximum brightness temperature of incoherent processes. Examples of known radio transient sources are indicated, including giant pulses and “nanogiant” pulses from the Crab pulsar (Hankins et al. 2003), radio pulses from XTE J1810–197 (Camilo et al. 2006) and other pulsars from the Australia Telescope National Facility pulsar catalogue (Manchester et al. 2005); the microquasar GRS 1915+105 (Mirabel & Rodríguez 1994); radio flares from V4641 Sgr (Hjellming et al. 2000), the brown dwarf LP 944–20 (Berger et al. 2001), and the magnetar SGR 1806–20 (Gaensler et al. 2005; Cameron et al. 2005); the Galactic centre radio transient J1745–3009 (Hyman et al. 2005); pulses from the ultracool dwarf TVLM 513–46546 (Hallinan et al. 2007); as well as radio emission from the Sun and Jupiter. Red lines indicate the expected sensitivity of ASKAP to sources at distances of 10 pc, 1 kpc and 1 Mpc.

within several hours.

Large-scale blind surveys for radio transients have been performed by a number of groups (e.g., Amy et al. 1989; Katz et al. 2003; Matsumura et al. 2007) over

the last few decades. These surveys have been severely limited in their sensitivity, sky coverage, and cadence. In spite of this, they have revealed that the radio sky contains many transient objects, the identities of some of which remain mysterious.

For example, Hyman et al. (2005) discovered a bursting transient towards the Galactic Centre which lasted for only a few minutes with a flux density in excess of 2 Jy at 330 MHz. Subsequent observations showed that the bursts repeated multiple times with a period of  $\sim 77$  minutes and a burst length of  $\sim 10$  minutes. The identification of this source remains unclear (e.g. Kaplan et al. 2008; Roy et al. 2010), although similarities with bursts from ultracool dwarfs (Hallinan et al. 2007) are intriguing. Another example is the well-known Lorimer burst, discovered in an archival search of Parkes pulsar data. Lorimer et al. (2007) discovered a single millisecond pulse of extragalactic origin with the astonishing peak flux density of 30 Jy at 1.4 GHz (also see Keane et al. 2011 for a second example). Although the nature of this source has been questioned, and recent work suggests the burst may be terrestrial (Burke-Spolaor et al. 2011), these new discoveries illustrate that the parameter space for radio variability has not been fully explored yet (Figure 2).

In this section we summarise existing blind and archival surveys for radio transients, focusing on ‘slow’ transients with variability on timescales of seconds or greater.

### 3.1 Transient Rates from Blind Surveys

Bower et al. (2007) presented the results from an archival survey of VLA data. Since the publication of this work there have been a number of blind radio transient surveys conducted on both archival data, and from the first datasets produced by next generation radio telescopes. Most of these have had the dual aim of (a) searching for individual objects of interest and trying to identify them; and (b) characterising the population statistics of radio transients in order to prepare for future surveys. Although the reported transient rates from Bower et al. (2007) have recently been revised, after a reanalysis of the data by Frail et al. (2012), the work established a framework for comparison of radio transient surveys.

Bower et al. (2007) proposed a metric for comparing the results from blind surveys for radio transients: the two-epoch equivalent snapshot rate, which is equivalent to the transient surface density. Figure 3 shows an updated version of their plot, incorporating a subset of the surveys listed in Table 3 and discussed in the rest of this section. For surveys that made no detection, an upper limit is calculated using a Poisson distribution:

$$P(n) = e^{-\rho A} \quad (2)$$

where  $P(n)$  is the confidence level (i.e.  $P(n) = 0.05$  for 95% confidence),  $\rho$  is the areal density of transients in a two-epoch survey, equivalent to the snapshot transient rate, and  $A$  is the equivalent solid angle, calculated by multiplying the number of images  $N_T$  by the field of

view  $\Omega$ . We can rearrange Equation 2 in terms of  $\rho$  for the purpose of placing theoretical constraints on planned surveys:

$$\rho = -\frac{\ln P(n)}{\Omega \times N_T}. \quad (3)$$

We can improve Equation 3 by taking into account not only the total number of observations in a given survey  $N_T$ , but the intra-observation cadence  $\tau$ , out of total observing time per observation  $T$ . Doing this results in:

$$\rho = -\frac{\ln P(n)}{\Omega \times N_T \times (T/\tau)}. \quad (4)$$

For example, on the smallest timescale, ASKAP images will be produced every  $\tau = 5$  seconds; the total observing time per observation (in one scenario) could be  $T = 1$  hour and the field could be observed daily for two years ( $N_T = 730$ ). A five second ASKAP image will have an RMS of 1.4 mJy beam $^{-1}$ . If subsequent data or images are stacked together (either via the uv or image plane) the RMS will improve to:

$$\sigma = \frac{1.4 \text{ mJy beam}^{-1}}{\sqrt{N_T \times (T/\tau)}} \quad (5)$$

Equations 4 and 5 have been used to generate survey predictions from the VAST survey parameters described in Section 4, and these predictions are shown in Figure 3. We also include the predicted rates of a subset of known radio transients (see Frail et al. 2012, for further discussion). One limitation of the plot in Figure 3 is that it collapses a number of dimensions into a single two-dimensional plot. For example, it does not represent the cadence of the surveys, or any trends in transient detection rate with observing frequency.

Surveys typically fall into one of two categories: many repeated observations of a single field (e.g. Carilli et al. 2003; Bower et al. 2007), or a large area, few-epoch surveys (e.g. Levinson et al. 2002; Gal-Yam et al. 2006). The two-epoch equivalent snapshot rate is a valid metric for quantising the effectiveness of a survey, provided the timescale of transient behaviour is less than the cadence of the observations but longer than a single integration. In this case, a survey with many epochs is equivalent to a larger area survey with two epochs. Further survey metrics have been developed to account for observation cadence and source evolution timescales as well (e.g., Cordes 2007).

### 3.2 Archival Surveys

The first large radio archival survey of variability and transient behaviour was conducted by Levinson et al. (2002) who covered a large fraction of the sky with two epochs by comparing the NVSS (Condon et al. 1998) and FIRST (White et al. 1997) surveys, both at 1.4 GHz. Various corrections were made to account for the different footprint of each survey and their differing resolutions (45'' for NVSS compared to 5'' for FIRST). Nine possible radio transients were reported by Levinson et al. (2002), but follow-up of these sources with the VLA by Gal-Yam et al. (2006) ruled out five of these as false candidates and two as



Table 3: Summary of snapshot rates for transient and variables radio sources reported in the literature. The results are organised according to upper limits based on non-detections (top section); transient detections (middle section); and detections of highly variable radio sources (bottom section).

Survey/Reference	$S_{Min}$ (mJy)	$S_{Max}$ (mJy)	Rate (deg <sup>-2</sup> )	Timescale	Frequency (GHz)	Epochs (N)
Bower et al. (2007)	>0.09	—	<6	1 year	4.8 & 8.4	17
Bower et al. (2007) & Frail et al. (2012) <sup>†</sup>	0.2	—	<3	2 months	4.8 & 8.4	96
Bower et al. (2007) & Frail et al. (2012) <sup>†</sup>	>0.37	—	< 6 × 10 <sup>-1</sup>	20 mins - 7 days	4.8 & 8.4	944
PiGSS-I/Bower et al. (2010)(A)*	>1	—	<1	1 month	3.1	75
PiGSS-II/Bower et al. (2011)(B)*	> 5	—	<0.18	1 month	3.1	78
FIRST-NVSS/Gal-Yam et al. (2006)	>6	—	<1.5×10 <sup>-3</sup>	days to months	1.4	2
Bell et al. (2011)	>8 (8 $\sigma$ )	—	<0.032	4.3 - 45.3 days	1.4, 4.8 & 8.4	5037
PiGSS-I/Bower et al. (2010)(A)*	>10	—	<0.3	1 month	3.1	75
PiGSS-II/Bower et al. (2011)(B)*	> 15	—	<0.025	1 day	3.1	78
ATATS - I/Croft et al. (2010)	>40	—	<0.004	81 days - ~ 15 years	1.4	12
Bower & Saul (2011)(A)*	>70	—	<3×10 <sup>-3</sup>	1 day	1.4	1852
ATATS - II/Croft et al. (2011)	>350	—	<6×10 <sup>-4</sup>	minutes to days	1.4	12
Bower & Saul (2011)(B)*	>3000	—	>9×10 <sup>-4</sup>	1 day	1.4	1852
Lazio et al. (2010)	>2.5×10 <sup>6</sup> (5 $\sigma$ )	—	<9.5×10 <sup>-8</sup>	5 mins	0.0738	~1272
JA12 <sup>T</sup> /Jaeger et al. (2012)	2.1	—	<0.12	1 day to 3 months	0.325	6
Ofek <sup>T</sup> /Ofek et al. (2011)	2.4	—	0.039	1 day–2 years	4.9	16
MOST <sup>T</sup> /Bannister et al. (2011a) <sup>‡</sup>	14(5 $\sigma$ )	6500	1.5 × 10 <sup>-3</sup>	days–years	0.843	3011
M09 <sup>T</sup> /Matsumura et al. (2009)	1000	4300	3×10 <sup>-3</sup>	~ 1 day	1.4	—
C03 <sup>V</sup> /Carilli et al. (2003)	>0.1	—	<18	19 days and 17 months	1.4	5
F03 <sup>V</sup> /Frail et al. (2003)	>0.25	—	5.8	~ 1 day	5 & 8.5	—
Becker <sup>V</sup> /Becker et al. (2010)	>1	—	1.6	~ 15 years	4.8	3
GT86 <sup>V</sup> /Gregory & Taylor (1986)	4.6	—	1×10 <sup>-3</sup>	1 day - 5 years	4.9	16
MOST <sup>V</sup> /Bannister et al. (2011a)	>14	—	0.268	days–years	0.843	3011

<sup>†</sup>Note that although Bower et al. (2007) report 10 detections, the reanalysis by Frail et al. (2012) results in only 4 (uncertain) detections. The snapshot rate calculated by Frail et al. (2012) assumes no detections, and we adopt this as the more conservative position for our analysis.

\*Bower et al. (2010), Bower et al. (2011) and Bower & Saul (2011) each state two different rates depending on flux density and timescale; we quote these separately as (A) and (B).

<sup>‡</sup>See also Bannister et al. (2011b).

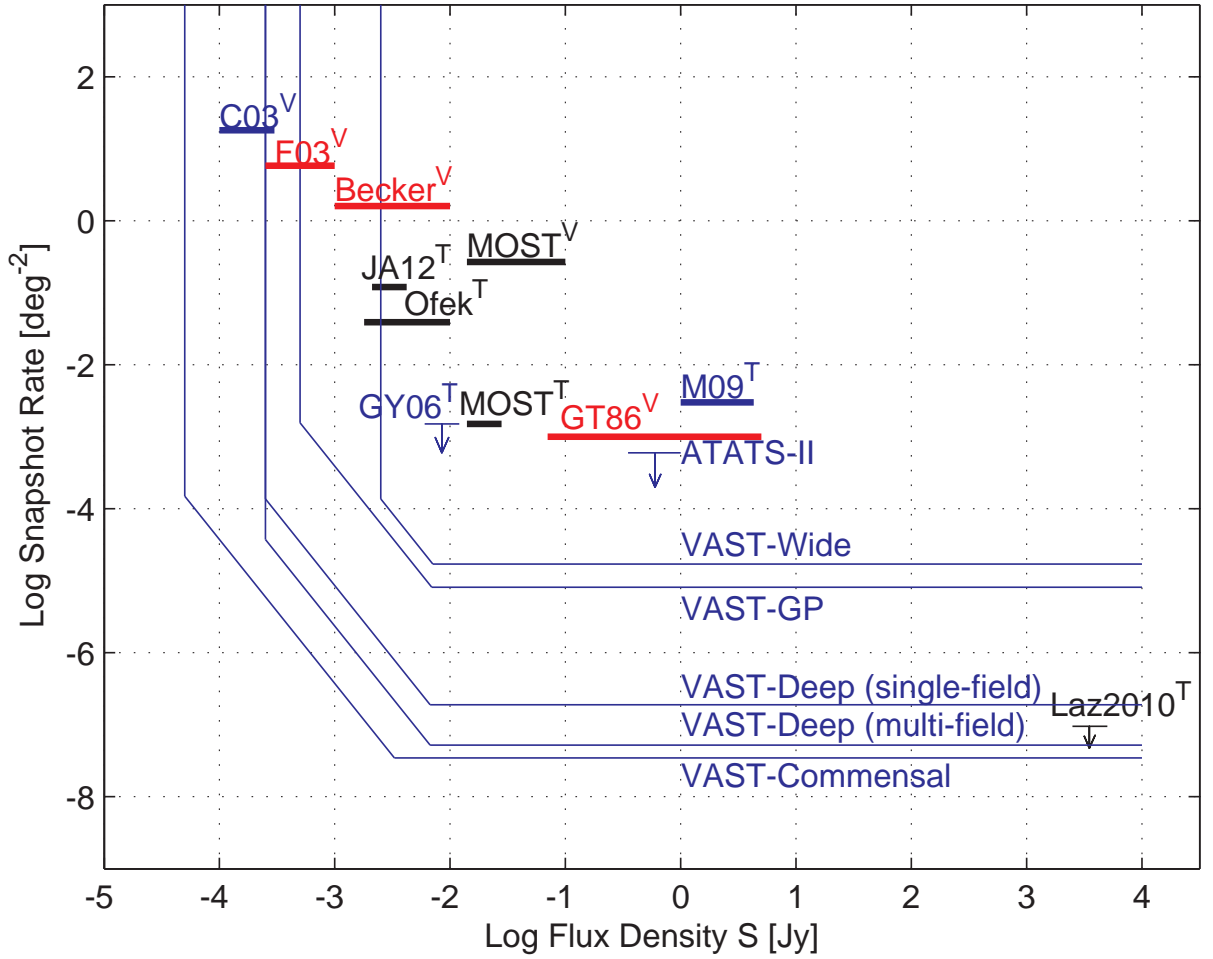


Figure 3: Log 2-epoch snapshot transient rate ( $\text{deg}^{-2}$ ) against log of the flux density (Jy) for surveys that report detections of transient and variables (thick lines), we also include a selection of surveys that report upper limits (thin lines with arrows) — see Table 3 for survey acronyms. The surveys are coloured according to frequency: black  $< 1$  GHz, blue =  $1 - 4.8$  GHz and red  $> 4.8$  GHz (see Table 3 for details). No corrections have been made for spectral index effects. Surveys labelled superscript ‘V’ denote detections of highly variable radio sources; those with superscript ‘T’ denote detections of transient type sources. The VAST survey predictions are indicated and organised by sub-survey. A description of each survey is given in Section 4. In each case the vertical segment denotes the RMS that can be achieved per observation. The horizontal segment indicates the upper limit that would be set if no transients or variables were detected in the entire survey. The joining line (between the horizontal and vertical) indicates the upper limits that could be placed by combining data from multiple five second snapshots up to the total available integration time per observation (see Table 3). See Section 3.1 for further discussion.

non-variable sources. The effective survey area was  $5990 \text{ deg}^2$ , hence with two detections the overall transient snapshot rate was  $\rho = 1.5 \times 10^{-3} \text{ deg}^{-2}$ .

Another large area survey was conducted by Bannister et al. (2011a), using archival data from the Molonglo Observatory Synthesis Telescope (Mills 1981) which operates at 843 MHz. This dataset had the advantage of a common resolution and comparable sensitivity in each epoch, but like the NVSS-FIRST comparison had a range of cadences between days and years for each field. Bannister et al. (2011a,b) detected 2 transient sources from a survey with a two-epoch equivalent

area of  $2776 \text{ deg}^2$ , giving an overall snapshot rate of  $\rho = 1.5 \times 10^{-3} \text{ deg}^{-2}$  for sources above 14 mJy.

An alternative approach has been to use archival calibrator fields which have many repeated observations but a relatively small field of view. Bell et al. (2011) analysed 5037 observations of VLA calibrator fields at three frequencies (1.4, 4.8 and 8.4 GHz). The observations ran from 1984 to 2008, with typical gaps between observations of between 4.3 and 45.3 days. No radio transients were detected and they therefore placed an upper limit of  $\rho < 0.032 \text{ deg}^{-2}$  for sources brighter than 8.0 mJy between 1.4 and 8.4 GHz.

Likewise Bower & Saul (2011) detected no transients in their analysis of archival observations of 3C 286 from the VLA. Their dataset spanned 23 years and included 1852 epochs at 1.4 GHz. They derive upper limits of  $\rho < 3 \times 10^{-3} \text{ deg}^{-2}$  at 70 mJy for timescales of  $\sim 1$  day and  $\rho < 9 \times 10^{-4} \text{ deg}^{-2}$  at 3 Jy for timescales of  $\sim 1$  minute. These results, combined with the Croft et al. (2010) results discussed in Section 3.3 below, currently provide the best limits on the number of slow, bright radio transients.

### 3.3 Gigahertz Surveys

There have been a number of blind imaging surveys specifically designed to search for transient and highly variable sources. Until recently these have generally been small area surveys with many epochs, due to the limited field of view of the telescopes used. Early surveys were motivated by searching for radio counterparts to GRBs and supernovae. For example, Frail et al. (1994) observed 15 epochs of a field centred on the location of GRB 940301 with the Dominion Radio Astrophysical Observatory Synthesis Telescope at 1.4 GHz. These observations were taken between 3 days and 99 days after the initial event and made no detections, which resulted in an upper limit of 3.5 mJy on any time-variable radio sources.

Carilli et al. (2003) observed the Lockman Hole region to a sensitivity limit of  $100 \mu\text{Jy}$  at 1.4 GHz with five epochs over a period of 17 months. They found that 2% of the sources within their sample were highly variable (defined as  $\Delta S \geq 50\%$ ) and hence derived a rate of  $\rho = 18 \text{ deg}^{-2}$ . Also in search of highly variable radio sources, Becker et al. (2010) surveyed the Galactic plane and found  $\rho = 1.6 \text{ deg}^{-2}$  above a sensitivity 1 mJy at 4.8 GHz. At low Galactic latitude and sampling a multitude of timescales (days to years) Ofek et al. (2011) used the VLA to study transient and variables in a field covering  $2.66 \text{ deg}^2$  to a sensitivity limit of  $\sim 100 \mu\text{Jy}$  at 5 GHz. One transient candidate was reported and a thorough discussion of the variable radio sources within the field was given.

Matsumura et al. (2009) summarise the detections of nine candidate transient sources from the Nasu 1.4 GHz wide-field survey (see Kuniyoshi et al. 2007; Matsumura et al. 2007; Niinuma et al. 2007; Kida et al. 2008; Niinuma et al. 2009, for full details). The sources were detected in a drift scan mode, and have flux densities greater than 1 Jy with typical timescales of minutes to days. (Note, however, the questions raised about these results by Croft et al. 2010, 2011 and the discussion about the resulting transient event rates by Ofek et al. 2011.)

As briefly discussed earlier in Section 3.1, Bower et al. (2007) reported the detection of ten transient radio sources in archival VLA data at 4.8 and 8.4 GHz. Frail et al. (2012) recently reanalysed this dataset and reported that more than half of these transients were either caused by rare data reduction artifacts, or that the detections had a lower S/N after re-reduction. The Bower et al. (2007) study initially predicted a rate of events over the whole sky of  $\rho \sim 1.5 \text{ deg}^{-2}$ , which is diminished to potentially an order of magnitude less, depending on how the lower S/N level transients are

interpreted. For Figure 3 we have adopted the conservative snapshot rate calculated by Frail et al. (2012), which assumes no detections. See Frail et al. (2012) for further discussion.

The Allen Telescope Array (ATA; Welch et al. 2009) is the first of the next generation telescopes to have been used for large area transient surveys. The Pi GHz Sky Survey (PiGSS) survey was carried out on the ATA at 3.1 GHz and covered  $10\,000 \text{ deg}^2$  of the sky with a minimum of two epochs for each field with an RMS sensitivity of  $\sim 1$  mJy (Bower et al. 2010).

The first data release (PiGSS-I) is a  $10 \text{ deg}^2$  region in the Boötes constellation, with 75 daily observations over a 4 month period. Bower et al. (2010) identify one object present in PiGSS that is not in archival radio catalogues (NVSS and WENSS). From this they set an upper limit on the snapshot rate of transients of  $\rho < 1 \text{ deg}^{-2}$  at 1 mJy sensitivity at 3.1 GHz. The second data release (PiGSS-II; Bower et al. 2011) focused on studying the daily and monthly transient and variable nature of the field. The upper limits of  $\rho < 0.025 \text{ deg}^{-2}$  at 15 mJy for one day timescales and  $\rho < 0.18 \text{ deg}^{-2}$  at 5 mJy for one month timescales were placed on the rates of transient events at 3.1 GHz.

The other major ATA survey is the The Allen Telescope Array Twenty-Centimeter Survey (ATATS; Croft et al. 2010), a  $690 \text{ deg}^2$  survey with 12 epochs of each field. Croft et al. (2010) presented results from comparing a single combined ATATS image with the NVSS. They did not detect any transients, placing an upper limit of  $\rho < 0.004 \text{ deg}^{-2}$  on the transient rate for sources with flux greater than 40 mJy at 1.4 GHz. Croft et al. (2011) extended this analysis presenting results for an individual epoch-to-epoch comparison for the ATATS dataset. They derived a limit of  $\rho < 6 \times 10^{-4} \text{ deg}^{-2}$  for transients brighter than 350 mJy at 1.4 GHz.

### 3.4 Low Frequency Surveys

Since many significant discoveries of radio transients have been made at low frequencies, the low frequency SKA pathfinders have enormous discovery potential.

The largest blind survey for slow transients at low frequencies was conducted by Lazio et al. (2010) with the Long Wavelength Demonstrator Array (LWDA). The LWDA is a 16 dipole phased array with all-sky imaging capabilities, operating at 73.8 MHz. A total of 106 hours of data (with a typical integration time of 5 minutes) was searched for transients. No detections were made above a flux density limit of 500 Jy. Lazio et al. (2010) quote an upper limit on the rate of  $1 \times 10^{-2} \text{ yr}^{-1} \text{ deg}^{-1}$  which converts to a rate of  $\rho < 9.5 \times 10^{-8} \text{ deg}^{-2}$  assuming 5 minute integrations.

In comparison, one of the deepest low frequency studies to date has been conducted by Jaeger et al. (2012). They used six VLA observations centred on the SWIRE Lockman Field at 325 MHz to study sources brighter than 0.2 mJy within a field of view of  $6.5 \text{ deg}^2$ . The timescales of observation ranged from 1 day to three months. One transient was reported and the authors argue that it is most likely coherent emission from a stellar flare. Using this one detection, Jaeger et al. (2012) set an upper limit on the snap-

shot rate of  $\rho < 0.12 \text{ deg}^{-2}$  for sources brighter than 2.1 mJy at 325 MHz.

Between the two comprehensive low frequency transient surveys discussed above (Lazio et al. 2010 and Jaeger et al. 2012) lies a large amount of parameter space both in frequency and time to explore. This will be made possible by two new instruments: the Murchison Widefield Array (MWA; Lonsdale et al. 2009; Tingay et al. 2012) and the Low Frequency Array (LOFAR; Heald et al. 2011), both currently being constructed and commissioned.

## 4 Survey Strategy

The detection and characterisation of radio transients requires both the sampling of new regions of parameter space, and the ability to monitor and follow up known or newly detected sources. As such, a successful transient instrument requires a large field-of-view, rapid survey capabilities, high sensitivity and dynamic range, as well as sufficient angular resolution for unambiguous source identification and follow-up. ASKAP will fulfill these criteria with its ability to achieve sub-millijansky sensitivity across the entire visible sky in a single day of observing.

Sources that vary on rapid timescales are necessarily compact. The  $10''$  resolution of ASKAP will permit source localisation to  $\sim$ arcsecond positional accuracy, enabling follow-up observations at other wave-bands.

The key to the efficient detection and monitoring of transient and variable sources is repeated observation of each field to sample a full range of cadences corresponding to the natural timescales of variability in the underlying source population. In addition to observing commensally with other ASKAP programs, we will run a campaign of complementary surveys designed to achieve our science goals. The VAST surveys build on the standard ASKAP imaging pipeline and thus do not include other possible modes of operation such as ‘Fly’s Eye’ monitoring (e.g., Siemion et al. 2012).

### 4.1 VAST Surveys

In this section we outline our strategy for each of the VAST surveys given in Table 2. We have based this survey strategy on an estimate of  $\sim 8500$  hours of observing time over the first five years of ASKAP operations, in addition to continuous commensal observing with other ASKAP projects. The survey parameters are summarised in Table 2. All of the dedicated (non-commensal) VAST surveys will be conducted with an observing frequency range of 1130–1430 MHz.

#### 4.1.1 VAST Wide

To detect rare events such as GRBs and SNe, and to map out the gas density distribution of the ISM across the whole sky, a large survey area is required. To fully sample the ESE light curves, a regular (preferably daily) cadence is desirable. Two years of data are required to discriminate between an annual cycle and intrinsic variability for IDV sources.

Hence our planned survey will cover 10 000 square degrees (comprised of 400 pointings, each of 30 square degrees). Integrating for 40 seconds per pointing will give an expected RMS sensitivity of  $0.5 \text{ mJy beam}^{-1}$ , with a total observing time (including overheads) of 6 hours per day. This survey will be repeated daily for 2 years, with a total time of 4380 hours. Figure 3 shows that it will be a substantial improvement over existing blind gigahertz surveys.

#### 4.1.2 VAST Deep

In our Deep survey we plan to observe 10 000 square degrees, with 400 pointings each observed for 1 hour. This would achieve an expected RMS sensitivity of  $50 \mu\text{Jy beam}^{-1}$ . Each of the fields will be observed by the EMU collaboration (see Section 4.2) for their continuum survey (Norris et al. 2011). While both the Deep and the Wide surveys observe the same amount of sky, they are complementary in cadence and sensitivity.

One field, coinciding with a field observed by the EMU collaboration, will be repeated daily for one year (Deep Single field). The other fields will be repeated 8 times at irregularly spaced intervals (Deep Multi-field). This is the optimal survey to detect as yet unknown source classes (see Bower et al. 2007) and detect GRBs and SNe to larger distances.

#### 4.1.3 VAST Galactic

A third component of VAST is a survey of 750 square degrees of the Galactic plane. This is a total of 30 pointings, plus an additional 3 pointings to cover the Large and Small Magellanic Clouds. Each pointing will be observed for 16 minutes, achieving an expected RMS sensitivity of  $0.1 \text{ mJy beam}^{-1}$ . The Galactic plane survey will require 9 hours of observing in a single day, repeated weekly for 1 year and supplemented by additional epochs at longer intervals.

## 4.2 Commensal Observing

In addition to the planned surveys, we will operate our transient detection pipeline commensally on all other ASKAP observations. Two projects, EMU (The Evolutionary Map of the Universe; Norris et al. 2011)<sup>5</sup> and WALLABY (The ASKAP HI All-Sky Survey; Koribalski, et al., in preparation)<sup>6</sup>, are particularly relevant as they will conduct extensive (possibly commensal) wide-field surveys:

**EMU:** a deep ( $10 \mu\text{Jy beam}^{-1}$  RMS) survey of the entire Southern Hemisphere, to a declination of  $\delta = +30^\circ$ .

**WALLABY:** a survey of neutral Hydrogen covering two thirds of the sky to a redshift of  $z = 0.5$ .

These projects plan to operate in ‘point and shoot’ mode, covering some region of the sky with significant integration time ( $\sim 12$  hours) per pointing and no revisits. This provides an excellent opportunity for

<sup>5</sup>See <http://www.askap.org/emu>.

<sup>6</sup>See <http://www.askap.org/wallaby>.



detection and monitoring of transient and variable behaviour on timescales of less than a day.

### 4.3 VAST Predictions

In Figure 3 we show the predictions parameterised via the method described in Section 3, using the VAST survey parameters given in Table 3. Note, the horizontal lines show the upper limits that would be placed if no transient or variable sources were detected. The vertical lines show the RMS that would be expected for the final integration time per image (see RMS sensitivity in Table 3). For the commensal survey the best RMS is found after six hours (not 12), assuming that two back-to-back six hour images are needed to search for transients, i.e., that there are no repeat observations. It can be seen from Figure 3 that the VAST surveys make orders of magnitude improvements on previous surveys.

VAST-Wide will use 4320 hours of dedicated observing time, which is the greatest of all the VAST sub-surveys (excluding commensal observing). It will observe 10 000 deg<sup>2</sup> for 40 seconds every day for two years, possibly longer. Observing for 40 seconds with ASKAP yields an RMS of  $\sim 0.5$  mJy beam<sup>-1</sup> and produces eight five-second images per observation. This is represented on Figure 3 and it casts VAST-Wide as less competitive than, for example, VAST-Deep (both multi-field and single field). The limits shown in Figure 3 for VAST-Wide however do not include upper limits on the rate and RMS calculations via combining data from multiple days observing.

For example, all the 40 second images taken daily for an entire week could be combined together and compared and searched for transients with the combined images from the next week, and so on. If technically feasible, such an approach will significantly increase the competitiveness of the VAST-Wide survey. Figure 3 also does not convey how many *known* radio sources will be monitored in each case. VAST-Wide will monitor all known sources in the Southern hemisphere above 2.5 mJy ( $5\sigma$ ) every day for two years, whereas VAST-Deep (one field) will monitor only a subset (30 deg<sup>2</sup>) to a greater sensitivity of 250  $\mu$ Jy ( $5\sigma$ ). We can approximate the ratio of sources that will be observed with VAST-Deep ( $N_{Deep}$ ) to VAST-Wide ( $N_{Wide}$ ) using the following expression (assuming a Euclidean source population):

$$\frac{N_{Deep}}{N_{Wide}} = \frac{A_{Deep}}{A_{Wide}} \times \left( \frac{\sigma_{Deep}}{\sigma_{Wide}} \right)^{-1.5} \simeq 10\% \quad (6)$$

where in this example  $A_{Deep}$  is the survey area of VAST-Deep (single field, 30 deg<sup>2</sup>);  $A_{Wide}$  the area of VAST-Wide (10,000 deg<sup>2</sup>).  $\sigma_{Deep}$  and  $\sigma_{Wide}$  represent the survey sensitivities respectively. VAST-Wide will therefore monitor large numbers of bright transients and variables, such as IDVs and AGN. VAST-Deep will be better optimised for studying the faint end of the transient and variable population.

## 5 Data Processing

The success of VAST hinges on our ability to detect transient sources on the timescales on which images are produced by ASKAP, at the same time as monitoring known variable sources that are in our field. Firstly, how to effectively deal with the large volumes of data, and secondly, how to find the ‘needles in a haystack’ — the most interesting transient sources. As discussed in Section 3, current blind surveys at 1.4 GHz suggest that less than 0.5% of sources observed by VAST will be transient or highly variable (Bannister et al. 2011a).

The operation of the ASKAP data pipeline is described in the Science Processing Document (Cornwell et al. 2012). The following sections address some of the issues most relevant to VAST, and discuss the design and operation of the VAST transient detection pipeline. The overall functionality of the VAST pipeline is illustrated in Figure 4 and described in Table 5. A discussion of the pipeline design is also provided by Banyer et al. (2011).

### 5.1 Data Rates

ASKAP is designed to be a real-time telescope, with most of the data processing done in online mode. Data from the phased array feed elements in each antenna will be sent to the beam former to create primary beams. The primary beams from each antenna are sent to the correlator to create visibility data. These data are then sent to the central processor for processing into science-ready data products: images and spectral cubes. The data rate from the correlator to the central processor is expected to be 2.5 GB/s, generating about 200 TB of raw data every 24 hours.

The central processor will operate three data processing pipelines in parallel: one pipeline will produce data cubes designed for spectral line analysis, with high frequency resolution; a second will produce data cubes designed for continuum surveys, with full polarisation and higher angular resolution, but reduced spectral resolution; the third pipeline is the transient detection pipeline which will produce data cubes with lower frequency resolution, but operating on 5-second timescales (see Table 4). Each 5-second image cube will be up to 16 GB (depending on the number of frequency channels generated and stored), which means that the transients pipeline will process around 12 TB of data an hour. Retaining some coarse frequency resolution will enable post-correlation rejection of radio frequency interference (see section 5.3 below) and also provide some minimal spectral information for detected transient sources.

### 5.2 Transient Detection Approach

Source detection is central to all upcoming radio transient surveys, and also those in other wave-bands such as the Large Synoptic Survey Telescope transients survey. In the image domain, transient source detection has typically followed one of two approaches: image subtraction or catalogue-based detection.

Table 4: VAST Image Products

Polarisations	1–4
Frequency channels	30
Field of view	7.5 deg <sup>2</sup>
Angular resolution	10''
Integration time	5 seconds
Image size (pixels)	4096 × 4096
Image cube size	up to 16 GB

Supernova search projects and other optical astronomy projects have typically used the image subtraction approach (for example, Riess et al. 2004; Alcock et al. 2000). In its simplest form, this involves subtracting the current image from some higher sensitivity master image. In practice, some corrections are required, for example each image can be convolved with a kernel to account for coordinate shifts and distortions (Alard & Lupton 1998; Alard 2000).

In radio images, the incomplete u-v coverage means that image subtraction often results in artefacts that can then be mistakenly identified as transient sources. A more robust approach is to firstly extract and measure sources in each image, and then search for transients by catalogue cross-matching and lightcurve analysis. The approach taken in the VAST transient detection pipeline is based on techniques described by Bannister et al. (2011a), developed for analysing the Molonglo Observatory Synthesis Telescope archive. This is also similar to the approaches taken in the Allen Telescope Array transients projects (Croft et al. 2010, 2011) and the LOFAR transient detection pipeline (Swinbank 2007).

We are also investigating the technique of transient detection in the visibility domain (Trott et al. 2011; Law & Bower 2012). These approaches have the advantage of dealing with well-behaved noise properties of visibilities and avoiding artefacts caused by the imaging and deconvolution processes. Trott et al. (2011) use statistical decision theory to derive detection limits of visibilities using models of the source and calibration effects of interferometers. Law & Bower (2012) present a computationally-simple, calibration-independent algorithm for transient detection based on the bispectrum, an interferometric closure quantity. By subtracting visibilities in time, the bispectrum can be used to quantify the significance of transients throughout the field of view. Visibility-based algorithms have great potential in the study of transients, because they take advantage of *a priori* information to reduce computational burden. This makes real-time processing easier, which may be a critical limitation of the large, computationally-demanding surveys described here. More work is needed to understand how these algorithms work on a large scale and whether they can be implemented for real-time analysis.

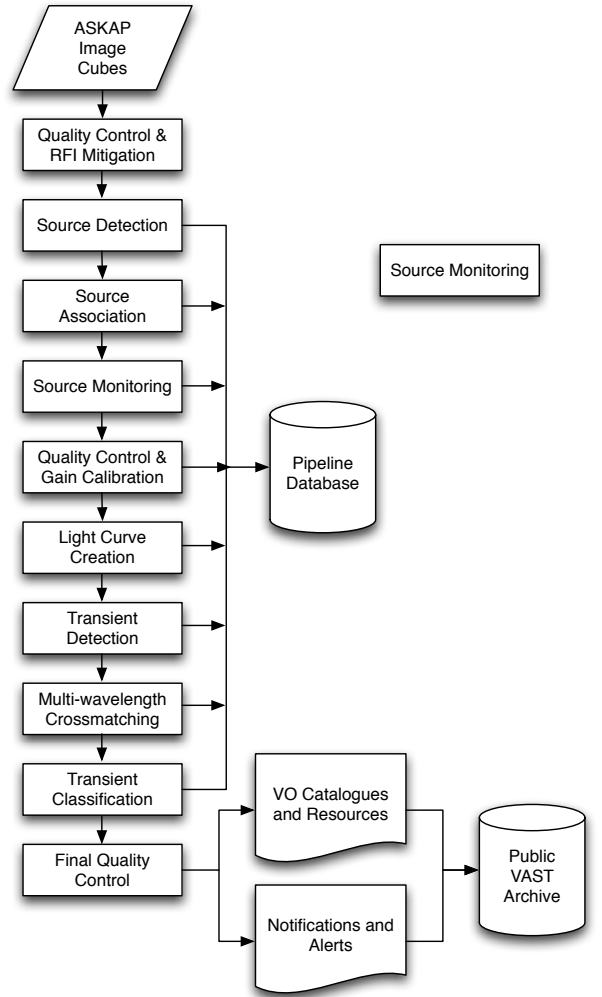


Figure 4: Illustration of the VAST pipeline functionality. A description of each stage is given in Table 5.

### 5.3 Rejection of Radio Frequency Interference

One of the most significant challenges for all radio telescopes is the impact of radio frequency interference (RFI) from man-made sources. Although the ASKAP site was chosen for its radio-quiet properties, signals from satellites used for communications, navigation, remote sensing and the military pervade all areas of the Earth. Similarly aircraft communication signals are present even in very remote areas. In addition, self-generated RFI is a challenge as faster digital equipment is used in the signal pathways. There is no universal method of RFI mitigation that is satisfactory for all types of observing techniques and RFI signals (Fridman & Baan 2001; Briggs & Kocz 2005). For the data rates expected with VAST, an automated approach with a combination of techniques will be essential.

RFI mitigation techniques range in sophistication

Table 5: Functionality of the VAST pipeline components, as illustrated in Figure 4.

Function	Description
Quality control & RFI mitigation	Initial quality control on the image cubes output from the ASKAP pipeline. Poor quality images will be flagged or ignored. Possible RFI mitigation strategies are discussed in Section 5.3.
Source detection	Blind source finding is performed on every image as discussed in Section 5.4. This produces a list of detected sources with position, flux, size and local RMS.
Source association	Each detected source is associated with one from the master list by finding the master source at the same position as the detection. If no master source is found at that position then a new master source is added.
Source monitoring	Master sources known to be in the field that were not detected by the source finder are measured by attempting to fit a Gaussian at the position where the source is expected to be. If no source can be fit then an upper limit on the flux at that position is taken instead.
Quality control & gain calibration	The image is assessed for quality by comparing the measured parameters of the detected sources to a list of well-characterised sources using the techniques described by Bannister et al. (2011a).
Light curve creation	All measurements for each source are collated into a radio light curve — a collection of flux measurements for all epochs where the source was observed.
Transient detection	The updated light curves are analysed to detect transient or variable behaviour. These sources will go on list for monitoring, and cross-matching with external catalogues.
Multi-wavelength cross-matching	Sources of interest are cross-matched with external catalogues to provide additional information on the likely causes of variability, and to collect extra information for classification and identification.
Transient classification	Sources that exhibit transient or variable behaviour are classified using a machine-learning algorithm. Possible approaches are discussed in Section 5.5.
Quality control & final data products	After a final quality control stage, sources that appear genuine with a high degree of confidence will be released to the community via VOEvent notifications as discussed in Section 5.6. These alerts, along with science quality catalogues of transient and variable sources will be made public through the VAST archive.

from simple blanking of corrupted data to adaptive cancellation based on coherent detection. Simple blanking works well for strong short duration RFI, but is completely ineffective in removing low-level time-variable RFI that could mimic transient sources. Blanking removes signal as well as interference and reduces the sensitivity achieved in a given time. There is a growing literature on more sophisticated ideas for RFI elimination. Some of the techniques that we are exploring for VAST include:

1. Post-detection flagging based on recognition of

data that are affected by interference. This has been the traditional approach in radio interferometry and is appropriate for moderately strong interference, but it does come with a loss of sensitivity.

2. Adaptive cancellation techniques such as those described in Mitchell et al. (2005); Briggs & Kocz (2005); Kesteven et al. (2010). These require a separate reference antenna or beam on the sky to sample the RFI. An algorithm is used that adapts filter characteristics to match the RFI

spectral footprint. This approach requires a high sensitivity copy of the interferer.

3. Spatial filtering techniques that use the location of the known interferer or the relative arrival times of the signals to identify and separate the RFI signal from the astronomical signal (e.g. Kocz et al. 2010; Leshem et al. 2000; Ellingson & Hampson 2002; Smolders & Hampson 2002). These are appropriate for stationary or slowly moving RFI sources and require less sensitivity to the interfering source than the adaptive cancellation techniques.
4. Adaptive nulling techniques, which employ the capacity of the array to place response nulls on the RFI source. These are appropriate for very strong interfering sources and for telescopes such as ASKAP with independently steerable beams.
5. Post-correlation filtering techniques use software algorithms to remove RFI from interferometric data by subtraction of a weak, persistent RFI signal (Athreya 2009) or by implementation of low-pass time and frequency filters to remove variable RFI (Offringa et al. 2012). The advantage of these techniques is that they generally have a lower sensitivity requirement for detecting the interference. However, radio interferometers may not be able to use this approach to mitigate some types of slowly varying RFI signals.

## 5.4 Source Detection

Finding and characterising sources in each of the observed images are critical first steps for transient detection in the image domain. False positives in source finding have the ability to confound the detection of transients, reducing the confidence with which a transient can be detected. This has been identified as a problem in previous surveys, with source-finding errors contributing to the false detection rate (see, for example, Bannister et al. 2011a; Croft et al. 2011; Frail et al. 2012).

The source finding packages that are currently available (for example the Miriad package `sfind` (Hopkins et al. 2002) and the AIPS package `VSAD` (Condon et al. 1998)) are all able to achieve a high degree of completeness and reliability, typically better than 95% on standard radio continuum images. These packages have been extensively used in large-scale radio surveys. However, even a small number of missed real sources, or false positives, can cause problems in the identification of variable and transient sources.

In recent surveys the best source finding results have been obtained by using the above mentioned source finding packages in conjunction with a pipeline of other programs (see, for example, Mauch et al. 2003; Murphy et al. 2007; Norris et al. 2006). The refinement of the source finding pipeline is often done after all observations have been completed and the idiosyncrasies of the data are well understood. A lot of time and effort is spent normalising the input images, and filtering the resulting catalogues. In many cases a subset of the final

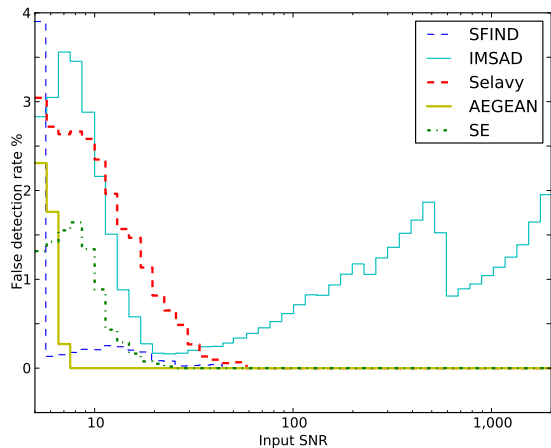


Figure 5: The false detection rate (FDR) for each of the source finding algorithms (from Hancock et al. 2012). No falsely detected sources are expected above an SNR of 5 for the area of sky simulated. The non-zero rate of false detections is due to poor source characterisation.

source catalogue is validated ‘by eye’ to ensure that the brighter sources are correctly represented. This approach produces high quality catalogues and is well suited to single epoch surveys, as is evidenced by the high level of completeness and reliability that these surveys are able to obtain.

For VAST, the volume of data will require a fully automated source finding process. The inability to archive every image that is observed for VAST will require that the source finding is done in real-time, and without the possibility of reprocessing. The large number of sources that are expected for a single day’s observation means that even with completeness and reliability of 99%, a source finding algorithm will be finding up to  $\sim 10,000$  false detections and missing up to  $\sim 10,000$  real sources each day. These false detections and missed sources far outweigh the number of true transient and variable sources that we expect to detect.

As part of the VAST design study process, Hancock et al. (2012) have undertaken a detailed analysis of some of the most widely used source finding packages in astronomy, including the newly developed Selavy (Whiting 2012) and Aegean (Hancock et al. 2012) algorithms. For compact sources, we find that the task of source *finding* (distinguishing between source and background pixels) is a solved problem that the various programs all achieve equally well. However, the task of source *characterisation* (measuring the various parameters of the found sources) is only accomplished well in the case of single isolated sources. For diffuse sources, the task of source finding is equally well achieved, however the task of source characterisation is not well defined. Hales et al. (2012) address some of the issues of extended source measurement. However, since we expect transient sources to be compact, this is not a



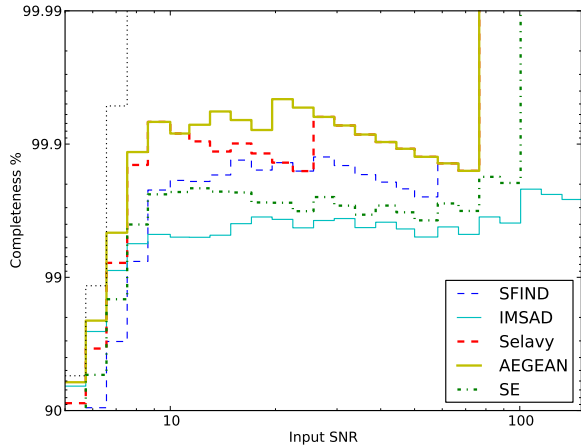


Figure 6: The completeness of each catalogue (from Hancock et al. 2012) as compared to the input catalogue. The coloured curves represent the completeness of the named source finders. The black dotted curve represents the expected completeness of an ideal source finder. The plateau in completeness above and SNR of 10 is due to sources that are poorly characterised.

significant issue for VAST.

Figures 5 and 6 show the false detection rate (FDR) and completeness of each of the source finding algorithms. The FDR and completeness were measured by injecting a known population of sources into a simulated image and processing this image with each of the different source finding algorithms. A comparison between the input and output catalogues resulted in: a set of sources common to both (real detections), a set of sources in the input but not the output (non-detections), and a set of sources that were in the output but not the input (false detections). The set of falsely detected sources consisted entirely of sources within an island of pixels that contains multiple components. These sources were so poorly characterised that they were not able to be linked to a source within the input catalogue. Sources that are poorly characterised also reduce the catalogue completeness. In Figures 5 and 6, the main reason that the source finding algorithms have non-zero FDR and less than ideal completeness is due to poor source *characterisation*, rather than poor source *finding*. In order to better characterise islands of pixels that contain more than a single component the Aegean program analyses the curvature of an image. As can be seen in Figures 5 and 6, Aegean is able to characterise islands of pixels that contain multiple components better than any of the other source finding approaches, typically achieving a completeness of over 99.9%. Further details of the source finding analysis and the operation of the Aegean source finding algorithm are presented by Hancock et al. (2012).

## 5.5 Classification of Transients

Given the large number of sources that we will likely detect with VAST, prompt classification will be critical for early identification and multi-wavelength follow-up. This is a challenging problem due to the sparsity and heterogeneity of the available data. There are two sets of information that can be used for this task: the light curve itself, and the existing archival information about that source or location on the sky.

Light-curve classification has been successfully used to classify variable stars at optical wavelengths (e.g. Richards et al. 2011, 2012). In order to explore the feasibility of light-curve classification for VAST, we used several approaches.

Firstly, we simulated light curves of five transient source types: SNe, RS CVn flare stars (fStar\_RSCVn), M-class flare stars (fStar\_dMe), XRBs, novae; two variable source types: IDVs and ESEs and a non-varying background (BG) for comparison (Lo et al. in preparation). We experimented with light-curve classification on the simulated light curves in two scenarios:

- *archival classification* where the classification is performed after a significant amount of light curve data has been collected; and
- *online classification*, which is classification of partial light curves in real-time for the purpose of triggering appropriate follow-up observations.

Secondly, we used light curves from the Palomar-Quest and Catalina Real-Time Transient Surveys to experiment further with a range of machine learning methods. The results of these experiments are discussed in the sections below.

### 5.5.1 Archival Classification

We approached the classification problem using the standard methodology of supervised machine learning with 10-fold cross validation. First we extracted a set of features from each light-curve. These include: time domain features which are simply the flux measurements at each time point (*tme*), statistical features (*stat*) such as the mean, standard deviation and skew, following Richards et al. (2011); statistical (*stat*) features built cumulatively (*stat-cum*) after observing a fixed number of observations in multiples of  $k$  ( $k$ ,  $2k$ , etc.); periodic features derived from the Lomb-Scargle periodogram (*lsp*); and wavelet features derived from the Haar wavelet (*wlet*). We then created two more feature sets that are concatenations of the above. Feature set *all-reps* is the concatenation of *stat-cum*, *lsp*, and *wlet*, and feature set *all* concatenates *tme*, *stat-cum*, *lsp*, and *wlet*. Details of all seven feature sets are discussed by Rebbapragada et al. (2012).

We used these features as input to a set of machine learning algorithms, which, when given a training set of examples, will learn the discriminative features for each source type. Using the WEKA toolkit (Hall et al. 2009), we experimented with three classification algorithms: Support Vector Machines (SVM; Cortes & Vapnik 1995), Random Forests (Breiman 2001) and the J48 Decision Tree (Quinlan 1986). Figure 7

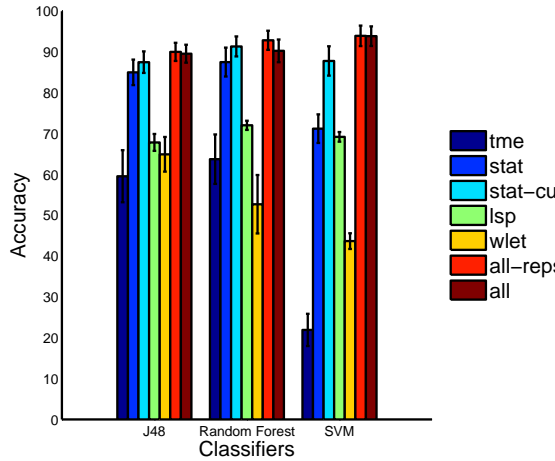


Figure 7: Comparison of overall classification performance (accuracy percentage) by feature set (*tme*, *stat*, *stat-cum*, *wlet*, *lsp*, *all-reps*, and *all*) and classification algorithm (J48, Random Forest, and SVM), with results grouped by classification algorithm.

shows the classification accuracy in the archival setting. The Random Forest is the best performing classifier when used with the combined feature sets, achieving 93% accuracy. This suggests archival classification using light curves alone is feasible with a reasonable degree of accuracy, given the large number of sources. However, the performance on real light curves is likely to be lower since our simulated light curves have ideal Gaussian noise and are aligned to the starting time of the event.

### 5.5.2 Online Classification

The aim of online classification is to identify sources of interest as soon as possible after new observations arrive. The earlier a classification is made, the more likely that a follow-up observation can be scheduled, and so online classification should run continuously as part of the real-time transient detection pipeline. We evaluated online classification using the first 30 days of simulated light curves. With so few observations, it is not possible to extract periodogram or wavelet features, hence we were limited to the statistical feature set and the time series values themselves.

Figure 8 shows the confusion matrix for our best performing algorithm: the J48 classifier built with *stat* features extracted from 30 days of observations. The overall classification accuracy achieved was 50% across eight source types. This is better than the baseline expected from random classification (12% accuracy) but probably not sufficient to be useful in scheduling follow-up observations. Accuracy was improved by grouping source types with similar light-curve characteristics together (grouping SNe and Novae, BG and ESE yielded an accuracy of 70%). Finally, accumulating more data led to higher accuracy ( $\sim 80\%$  accuracy with 100 days of observations).

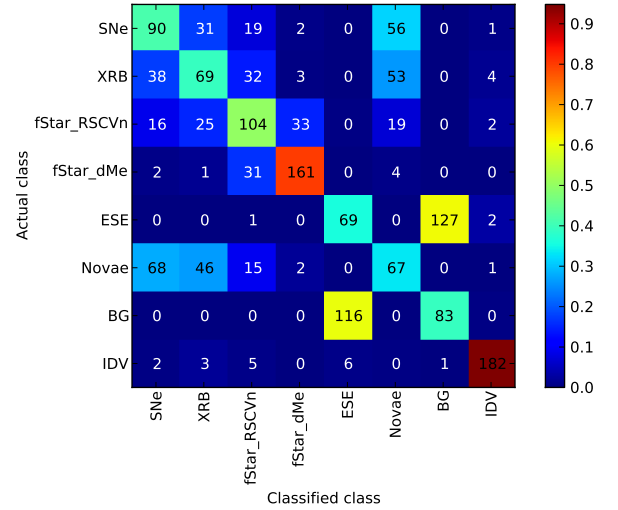


Figure 8: Online classification confusion matrix shown as a heat map (red = higher accuracy, blue = lower accuracy). The x-axis is the classified output class and the y-axis is the actual input class. We used 200 sources of each type and the number in each cell represents how these sources are classified. Overall accuracy at 30 days is 50%.

### 5.5.3 Other Machine Learning Approaches

Considerable insight into the classification problem can be gained from the current efforts on classification of transients in optical synoptic sky surveys (Bloom et al. 2008; Donalek et al. 2008; Mahabal et al. 2008a,b, 2010, 2011; Djorgovski et al. 2011). We experimented with several approaches, using the data from the Palomar-Quest and Catalina Real-Time Transient Surveys, as well as selected data sets from the literature.

Similar to the work described in Section 5.5.1 we parameterised observed optical light curves as feature vectors, and applied various machine learning techniques (e.g., decision trees) to their classification. A set of decision trees were trained using the feature vectors for various combinations of classes. To reduce the dimensionality of the input space, we applied a forward feature selection strategy that consisted of selecting a subset of features from the training set that best predict the test data by sequentially selecting features until there is no improvement in prediction. Each tree was built using the Gini diversity index as criterion for choosing the split; the splitting stopped when there is no further gain that can be made. To avoid over-fitting we used a 10-fold cross validation approach. The decision trees were pruned in order to choose the simplest one within one standard error of the minimum. When tested on a data set consisting of light curves of blazars, cataclysmic variables, and RR Lyrae stars, we achieved completeness in the range of  $\sim 83\%$  -  $97\%$ , and contamination in the range of  $\sim 4\%$  to  $13\%$  (Donalek et al., in prep.). In future work we will apply this to the classification of radio light curves.

Another approach we experimented with used two-dimensional distributions of magnitude changes for dif-

ferent time baselines for all possible epoch pairs in the data set (Djorgovski et al. 2011, Moghaddam et al., in preparation). These two-dimensional ( $\Delta m, \Delta t$ ) histograms can be viewed as probabilistic structure functions for the light curves of different types. Template distributions for different kinds of transients and variables are constructed using the reliably classified data with the same survey cadences, S/N, etc. For any newly detected variable or a transient, corresponding ( $\Delta m, \Delta t$ ) histograms are accumulated as the new data arrive, and we used a variety of metrics to compute the effective probabilistic distances from different templates. Initial testing indicates that we can obtain a high classification accuracy using this approach. We plan to generalise it to include triplets or even higher order sets of data points for multi-dimensional histograms.

One important lesson so far is that existing archival and contextual data will play a critical role in classification. This includes a spatial context (i.e., what is near the observed event on the sky — a possible host galaxy, a cluster, a SN remnant, etc.), the multi-wavelength context (has it been detected at other wavelengths, what is its broad spectral energy distribution), and the temporal context (what was its flux variability or a detection history in the previously obtained data). Some of this information can be readily extracted from the archives (e.g., flux measurements from different wavelengths, light curves at that location), but some - spatial context in particular - require a human judgment, e.g., is the apparent proximity to other objects (galaxies, clusters, etc.) likely to be relevant, and if so, what does it imply about the transient? Human inspection of vast numbers of transients does not scale to the massive data streams such as those contemplated here. We are currently experimenting with crowdsourcing approaches to harvesting of the relevant human pattern recognition skills and domain expertise, and their translation to machine-processable algorithms.

## 5.6 Alerts and Follow-up

ASKAP policy specifies that all data and results must be released as soon as they have passed scientific quality control. This is particularly important for transient detections, in which multi-wavelength follow-up is critical. VAST will generate publicly available data products on the following timescales:

**Seconds:** Triggers for individual transient events will be released through services such as VOEventNet<sup>7</sup> as soon as they are detected. These events will have quality flags representing our confidence about the detection and will allow immediate follow-up of new discoveries by the wider community.

**Minutes:** Triggers with value-added data such as classifications based on multi-wavelength archival crossmatching will be released as soon as our source classification pipeline has run.

**Days:** Measurements of every source in every field will be entered into our lightcurve database in near real-time. This will be available immediately to the VAST collaboration. Once quality control has been completed, updates to the database will be made publicly available.

**Months:** Scientific quality catalogues will be constructed from our lightcurve database. The results will be published and data made public as soon as possible after analysis is complete.

The quality control process for alerts will consist of a multi-tiered system in which there is automatic filtering of a large fraction of false detections followed by manual inspection of the remaining sources. This approach is taken by several existing transients surveys, for example in the Palomar Transients Factory (Gal-Yam et al. 2011). The detailed implementation of this general approach will be developed and refined during commissioning.

In addition, all images from VAST will be made available through the standard ASKAP archive. Where possible the VAST pipeline will use standard Virtual Observatory protocols such as VOEvents<sup>8</sup>.

### 5.6.1 Very Long Baseline Interferometry

As noted earlier, radio sources that are highly variable on short timescales are necessarily intrinsically compact. ASKAP's angular resolution of  $\sim 10''$  will be high enough to localise variable sources on the sky and measure variability, but will be orders of magnitude too low to resolve the evolving structures that give rise to the variations in integrated flux. In order to resolve the structure and evolution of the radio source, much higher angular resolution imaging is required, provided by Very Long Baseline Interferometry (VLBI).

Alerts from the VAST pipelines will be used to trigger observations using the Australian Long Baseline Array (LBA), now including telescopes from Western Australia (ASKAP) to New Zealand (Warkworth), allowing milliarcsecond-scale imaging. Due to the very large number of VAST alerts, targets for VLBI follow-up will have to be chosen selectively and will need to be restricted to the most interesting or rare classes of transient and variable sources. The Australian LBA is very well suited to this type of follow-up, as evidenced by previous programs that have taken triggers for transient radio sources from MOST or the ATCA (Tingay et al. 1995; Miller-Jones et al. 2012).

The current LBA has a sensitivity of  $60\mu\text{Jy beam}^{-1}$  for a one hour integration at 2.3 and 8 GHz. This is comparable to the sensitivity of ASKAP at 1.4 GHz. With ASKAP incorporated as an element of the LBA, the sensitivity at 1.6 GHz would be  $40\mu\text{Jy beam}^{-1}$  for a one hour integration. Hence a 12 hour observation with the LBA would be able to detect any unresolved source found by the VAST surveys.

The recent advent of real-time eVLBI on the LBA means that the VLBI observations can be rapidly processed and analysed. Thus subsequent observations

<sup>7</sup>See <http://voeventnet.caltech.edu>.

<sup>8</sup>See <http://www.ivoa.net/Documents/VOEvent>.

can be adapted based on the initial results. This facility will benefit studies of new transients that will be detected by ASKAP.

## 6 Conclusions

Next generation radio telescopes such as the EVLA, LOFAR, MWA and ASKAP will have a substantial impact on our ability to discover and monitor radio transient and variable sources. The VAST survey will play an important role in the study of a range of phenomena, from GRB afterglows and supernovae through to cataclysmic variables and radio flare stars, as well as propagation effects in the interstellar medium.

The three VAST surveys (VAST-Wide, VAST-Deep and VAST-Galactic) will allow us to do a comprehensive search of radio transient parameter space. It is likely that in addition to discovering a large number of known classes of objects, we will also discover new types of transient and variable phenomena. Our automatic transient detection pipeline will be able to run simultaneously with other major ASKAP observing projects, meaning that we will also analyse data in commensal mode.

All data products produced by VAST will be released to the wider community as soon as possible after passing quality control standards. Wherever possible this will be done through the use of Virtual Observatory tools and protocols.

In this paper we have summarised the science goals of VAST, described our current results in algorithm development for source-finding and transient light curve classification, and made predictions about the likely source populations that will be observed by VAST.

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## References

- Abadie, J., et al. 2010, *Classical and Quantum Gravity*, 27, 173001
- Abbasi, R., et al. 2012, *Nature*, 484, 351
- Abdo, A. A., et al. 2010, *ApJ*, 722, 520
- Accadia, T., et al. 2011, *Classical and Quantum Gravity*, 28, 114002
- Alard, C. 2000, *A&AS*, 144, 363
- Alard, C., & Lupton, R. H. 1998, *ApJ*, 503, 325
- Alcock, C., et al. 2000, *ApJ*, 541, 734
- Amy, S. W., et al. 1989, *PASA*, 8, 172
- Archibald, A. M., Kaspi, V. M., Bogdanov, S., Hessels, J. W. T., Stairs, I. H., Ransom, S. M., & McLaughlin, M. A. 2010, *ApJ*, 722, 88
- Archibald, A. M., et al. 2009, *Science*, 324, 1411
- Armstrong, J. W., Rickett, B. J., & Spangler, S. R. 1995, *ApJ*, 443, 209
- Armstrong, J. W., Spangler, S. R., & Hardee, P. E. 1977, *AJ*, 82, 785
- Athreya, R. 2009, *ApJ*, 696, 885
- Ball, L. 2009, ASKAP Survey Science Projects, Tech. rep., CSIRO Australia Telescope National Facility
- Bannister, K. W., Murphy, T., Gaensler, B. M., Hunstead, R. W., & Chatterjee, S. 2011a, *MNRAS*, 412, 634
- . 2011b, *MNRAS*, 418, 2813
- Banyer, J., et al. 2011, in *ADASS XXI* (arXiv 1201.3130), ASP Conference Series
- Barnacka, A., Glicenstein, J.-F., & Moudden, Y. 2011, *A&A*, 528, L3
- Bastian, T. S., Dulk, G. A., & Chanmugam, G. 1988, *ApJ*, 324, 431
- Becker, R. H., Helfand, D. J., White, R. L., & Proctor, D. D. 2010, *AJ*, 140, 157
- Bell, M. E., et al. 2011, *MNRAS*, 415, 2



- Berger, E. 2006, *ApJ*, 648, 629
- Berger, E., Zauderer, A., Pooley, G. G., Soderberg, A. M., Sari, R., Brunthaler, A., & Bietenholz, M. F. 2012, *ApJ*, 748, 36
- Berger, E., et al. 2001, *Nature*, 410, 338
- Bignall, H., et al. 2003, *ApJ*, 585, 653
- Bignall, H. E., Macquart, J.-P., Jauncey, D. L., Lovell, J. E. J., Tzioumis, A. K., & Kedziora-Chudczer, L. 2006, *ApJ*, 652, 1050
- Blandford, R. D., & Payne, D. G. 1982, *MNRAS*, 199, 883
- Bloom, J. S., Starr, D. L., Butler, N. R., Nugent, P., Rischard, M., Eads, D., & Poznanski, D. 2008, *Astronomische Nachrichten*, 329, 284
- Bloom, J. S., et al. 2011, *Science*, 333, 203
- Bond, H. E., White, R. L., Becker, R. H., & O'Brien, M. S. 2002, *PASP*, 114, 1359
- Bookbinder, J. A., & Lamb, D. Q. 1987, *ApJ*, 323, L131
- Bower, G. C., & Saul, D. 2011, *ApJ*, 728, L14
- Bower, G. C., Whysong, D., Blair, S., Croft, S., Keating, G., Law, C., Williams, P. K. G., & Wright, M. C. H. 2011, *ApJ*, 739, 76
- Bower, G. C., et al. 2007, *ApJ*, 666, 346
- . 2010, *ApJ*, 725, 1792
- Breiman, L. 2001, *Machine Learning*, 45, 5
- Briggs, F. H., & Kocz, J. 2005, *Radio Science*, 40, 5
- Brunthaler, A., Menten, K. M., Reid, M. J., Henkel, C., Bower, G. C., & Falcke, H. 2009, *A&A*, 499, L17
- Brunthaler, A., et al. 2010, in *Proceedings of the 10th European VLBI Network Symposium and EVN Users Meeting: VLBI and the new generation of radio arrays*
- Burke-Spolaor, S. 2010, *MNRAS*, 1574
- Burke-Spolaor, S., Bailes, M., Ekers, R., Macquart, J.-P., & Crawford, F. 2011, *ApJ*, 727, 18
- Cameron, P. B., et al. 2005, *Nature*, 434, 1112
- Camilo, F., et al. 2006, *Nature*, 442, 892
- Carilli, C. L., Ivison, R. J., & Frail, D. A. 2003, *ApJ*, 590, 192
- Cenko, S. B., et al. 2012a, *MNRAS*, 420, 2684
- . 2012b, *ApJ*, 753, 77
- Chakraborti, S., Ray, A., Soderberg, A. M., Loeb, A., & Chandra, P. 2011, *Nature Communications*, 2
- Chandra, P., & Frail, D. A. 2012, *ApJ*, 746, 156
- Chippendale, A. P., O'Sullivan, J., Reynolds, J., Gough, R., Hayman, D., & Hay, S. 2010, in *IEEE International Symposium on Phased Array Systems and Technology*, 648–652
- Chomiuk, L., et al. 2012, *ApJ*, 750, 164
- Clegg, A. W., Fey, A. L., & Fiedler, R. L. 1996, *ApJ*, 457, L23
- Clegg, A. W., Fey, A. L., & Lazio, T. J. W. 1998, *ApJ*, 496, 253
- Condon, J. J., Cotton, W. D., Greisen, E. W., Yin, Q. F., Perley, R. A., Taylor, G. B., & Broderick, J. J. 1998, *AJ*, 115, 1693
- Cordes, J. M. 2007, *The SKA as a Radio Synoptic Survey Telescope: Widefield Surveys for Transients, Pulsars and ETI*, Tech. rep., SKA Memo 97
- Cordes, J. M. 2012, in *IAU Symposium, Vol. 285, IAU Symposium*, 49–54
- Cordes, J. M., & Wolszczan, A. 1986, *ApJ*, 307, 27
- Cordes, J. M., et al. 2004, *New Astronomy Reviews*, 48, 1459
- Cornwell, T., et al. 2012, *ASKAP Science Processing Document*, Tech. rep., CSIRO Astronomy and Space Science
- Cortes, C., & Vapnik, V. 1995, *Machine Learning*, 20, 273
- Coward, D., et al. 2012, *MNRAS*, submitted (arXiv 1202.2179)
- Croft, S., Bower, G. C., Keating, G., Law, C., Whysong, D., Williams, P. K. G., & Wright, M. 2011, *ApJ*, 731, 34
- Croft, S., et al. 2010, *ApJ*, 719, 45
- Dalal, N., Griest, K., & Pruet, J. 2002, *ApJ*, 564, 209
- Davis, S. W., Narayan, R., Zhu, Y., Barret, D., Farrell, S. A., Godet, O., Servillat, M., & Webb, N. A. 2011, *ApJ*, 734, 111
- De Jager, O. C., Meintjes, P. J., O'Donoghue, D., & Robinson, E. L. 1994, *MNRAS*, 267, 577
- Dennett-Thorpe, J., & de Bruyn, A. G. 2002, *Nature*, 415, 57
- Dennison, B., & Condon, J. J. 1981, *ApJ*, 246, 835
- Dennison, B., et al. 1987, *ApJ*, 313, 141
- Dewdney, P. E., Hall, P. J., Schilizzi, R. T., & Lazio, T. J. L. W. 2009, *IEEE Proceedings*, 97, 1482
- Djorgovski, S. G., et al. 2011, in *Statistical Analysis and Data Mining*, in press

- Donalek, C., Mahabal, A., Djorgovski, S. G., Marney, S., Drake, A., Glikman, E., Graham, M. J., & Williams, R. 2008, in *American Institute of Physics Conference Series*, Vol. 1082, American Institute of Physics Conference Series, ed. C. A. L. Bailer-Jones, 252
- Draine, B. 1998, *ApJ*, 509, L41
- Dulk, H. E., Bastian, T. S., & Chanmugam, G. 1983, *ApJ*, 273, 249
- Ebisuzaki, T., et al. 2001, *ApJ*, 562, L19
- Ellingson, S. W., & Hampson, G. A. 2002, *IEEE Transactions on Antennas and Propagation*, 50, 25
- Esquej, P., Saxton, R. D., Freyberg, M. J., Read, A. M., Altieri, B., Sanchez-Portal, M., & Hasinger, G. 2007, *A&A*, 462, L49
- Esquej, P., et al. 2008, *A&A*, 489, 543
- Fairhurst, S. 2011, *Classical and Quantum Gravity*, 28, 105021
- Falcke, H., K rding, E., & Markoff, S. 2004, *A&A*, 414, 895
- Farrell, S. A., Webb, N. A., Barret, D., Godet, O., & Rodrigues, J. M. 2009, *Nature*, 460, 73
- Fassnacht, C. D., Pearson, T. J., Readhead, A. C. S., Browne, I. W. A., Koopmans, L. V. E., Myers, S. T., & Wilkinson, P. N. 1999, *ApJ*, 527, 498
- Fender, R. 2006, *Jets from X-ray binaries*, ed. W. H. G. Lewin & M. van der Klis, 381–419
- Fender, R., et al. 2004, *MNRAS*, 355, 1105
- Feng, H., & Soria, R. 2011, *New Astronomy Reviews*, 55, 166
- Fiedler, R., Pauls, T., Johnston, K. J., & Dennison, B. 1994a, *ApJ*, 430, 595
- Fiedler, R. L., Johnston, K. J., Waltman, E. B., & Ghigo, F. 1992, *IAU Circular*, 5527, 2
- Fiedler, R. L., et al. 1987a, *ApJS*, 65, 319
- 1987b, *Nature*, 326, 675
- 1994b, *ApJ*, 430, 581
- Fielder, R., Dennison, B., Johnston, K. J., Waltman, E. B., & Simon, R. S. 1994, *ApJ*, 430, 581
- Fong, W.-f., et al. 2012, *ApJ*, submitted (arXiv 1204.5475)
- Fornasa, M., & Bertone, G. 2008, *International Journal of Modern Physics D*, 17, 1125
- Frail, D., et al. 2005, *ApJ*, 619, 994
- Frail, D. A., Kulkarni, S. R., Berger, E., & Wieringa, M. H. 2003, *AJ*, 125, 2299
- Frail, D. A., Kulkarni, S. R., Ofek, E. O., Bower, G. C., & Nakar, E. 2012, *ApJ*, 747, 70
- Frail, D. A., et al. 1994, *ApJ*, 437, L43
- Fridman, P. A., & Baan, W. A. 2001, *A&A*, 378, 327
- Gaensler, B. M., et al. 2005, *Nature*, 434, 1104
- Gal-Yam, A., et al. 2006, *ApJ*, 639, 331
- 2011, *ApJ*, 736, 159
- Gallo, E., Fender, R. P., & Pooley, G. G. 2003, *MNRAS*, 344, 60
- Gezari, S., et al. 2006, *ApJ*, 653, L25
- 2008, *ApJ*, 676, 944
- 2009, *ApJ*, 698, 1367
- 2012, *Nature*
- Giannios, D., & Metzger, B. D. 2011, *MNRAS*, 416, 2102
- Godet, O., et al. 2012, *ApJ*, 752, 34
- Gregory, P. C., & Taylor, A. R. 1986, *AJ*, 92, 371
- Grupe, D., Thomas, H.-C., & Leighly, K. M. 1999, *A&A*, 350, L31
- Gunn, A. G., Spencer, R. E., Abdul Aziz, H., Doyle, J. G., Davis, R. J., & Pavelin, P. E. 1994, *A&A*, 291, 847
- Hales, C. A., Murphy, T., Curran, J. R., Middelberg, E., Gaensler, B. M., & Norris, R. P. 2012, *MNRAS*, in press (arXiv 1205.5313)
- Hall, M., Frank, E., Holmes, G., Pfahringer, B., Reutemann, P., & Witten, I. H. 2009, *SIGKDD Explorations*, 11
- Hallinan, G., et al. 2007, *ApJ*, 663, L25
- Halpern, J. P., et al. 2005, *ApJ*, 632, L29
- Hancock, P., Gaensler, B. M., & Murphy, T. 2011, *ApJ*, 735, L35
- Hancock, P. J., Murphy, T., Gaensler, B. M., Hopkins, A., & Curran, J. R. 2012, *MNRAS*, 422, 1812
- Hankins, T. H., Kern, J. S., Weatherall, J. C., & Eilek, J. A. 2003, *Nature*, 422, 141
- Harry, G. M., et al. 2010, *Classical and Quantum Gravity*, 27, 084006
- Heald, G., et al. 2011, *Journal of Astrophysics and Astronomy*, 32, 589
- Heeschen, D. S. 1984, *AJ*, 89, 1111
- Hjellming, R. M. 1990, in *Lecture Notes in Physics*, Vol. 369, *Physics of classical novae*. IAU Coll. 122, ed. A. Cassatella & R. V. Viotti (Springer-Verlag: Berlin), 167–178

- Hjellming, R. M., et al. 2000, *ApJ*, 544, 977
- Hopkins, A. M., Miller, C. J., Connolly, A. J., Genovese, C., Nichol, R. C., & Wasserman, L. 2002, *AJ*, 123, 1086
- Hunstead, R. W. 1972, *Astrophysical Letters*, 12, 193
- Hyman, S. D., et al. 2005, *Nature*, 434, 50
- Jaeger, T. R., Hyman, S. D., Kassim, N. E., & Lazio, T. J. W. 2012, *AJ*, 143, 96
- Jenkins, E. B., & Tripp, T. M. 2001, *ApJS*, 137, 297
- Johnston, S., et al. 2007, *PASA*, 24, 174
- . 2008, *Experimental Astronomy*, 22, 151
- Kankare, E., et al. 2008, *ApJ*, 689, L97
- Kaplan, D. L., Hyman, S. D., Roy, S., Bandyopadhyay, R. M., Chakrabarty, D., Kassim, N. E., Lazio, T. J. W., & Ray, P. S. 2008, *ApJ*, 687, 262
- Kaplan, D. L., O’Shaughnessy, R., Sesana, A., & Volonteri, M. 2011, *ApJ*, 734, L37
- Katz, C. A., et al. 2003, *PASP*, 115, 675
- Keane, E. F., Kramer, M., Lyne, A. G., Stappers, B. W., & McLaughlin, M. A. 2011, *MNRAS*, 415, 3065
- Kedziora-Chudczer, L., et al. 1997, *ApJ*, 490, L9
- Kesteven, M., Manchester, R., Brown, A., & Hampson, G. 2010, in *RFI Mitigation Workshop*
- Kida, S., et al. 2008, *New Astronomy*, 13, 519
- Kocz, J., Briggs, F. H., & Reynolds, J. 2010, *AJ*, 140, 2086
- Komossa, S. 2006, *Mem. Soc. Ast. It.*, 77, 733
- Komossa, S., & Bade, N. 1999, *A&A*, 343, 775
- Komossa, S., & Greiner, J. 1999, *A&A*, 349, L45
- Körding, E., Falcke, H., & Corbel, S. 2006, *A&A*, 456, 439
- Körding, E., Rupen, M., Knigge, C., Fender, R., Dhawan, V., Templeton, M., & Muxlow, T. 2008, *Science*, 320, 1318
- Krolik, J. H., & Piran, T. 2012, *ApJ*, 749, 92
- Kulkarni, S. R., & Heiles, C. 1988, *Neutral hydrogen and the diffuse interstellar medium*, ed. K. I. Kellermann & G. L. Verschuur, 95–153
- Kuniyoshi, M., et al. 2007, *PASP*, 119, 122
- Law, C. J., & Bower, G. C. 2012, *ApJ*, 749, 143
- Lazio, T. J. W., Gaume, R. A., Claussen, M. J., Fey, A. L., Fiedler, R. L., & Johnston, K. J. 2001a, *ApJ*, 546, 267
- Lazio, T. J. W., Waltman, E. B., Ghigo, F. D., Fiedler, R. L., Foster, R. S., & Johnston, K. J. 2001b, *ApJS*, 136, 265
- Lazio, T. J. W., et al. 2000, *ApJ*, 534, 706
- . 2010, *AJ*, 140, 1995
- Leshem, A., van der Veen, A.-J., & Boonstra, A.-J. 2000, *ApJS*, 131, 355
- Levan, A. J., et al. 2011, *Science*, 333, 199
- Levinson, A., Ofek, E. O., Waxman, E., & Gal-Yam, A. 2002, *ApJ*, 576, 923
- Lien, A., Chakraborty, N., Fields, B. D., & Kembell, A. 2011, *ApJ*, 740, 23
- Lin, D., Carrasco, E. R., Grupe, D., Webb, N. A., Barret, D., & Farrell, S. A. 2011, *ApJ*, 738, 52
- Lo, K. K., et al. 2012, *MNRAS*, 421, 3316
- Lonsdale, C., et al. 2009, *Proceedings of the IEEE*
- Lorimer, D. R., et al. 2007, *Science*, 318, 777
- Lovell, J. E. J., et al. 2003, *AJ*, 126, 1699
- . 2008, *ApJ*, 689, 108
- Lyutikov, M. 2011, *Phys. Rev. D*, 83, 064001
- Macquart, J., et al. 2010, *PASA*, 27, 272
- Macquart, J.-P., & de Bruyn, A. G. 2007, *MNRAS*, 380, L20
- Macquart, J.-P., Kedziora-Chudczer, L., Rayner, D. P., & Jauncey, D. L. 2000, *ApJ*, 538, 623
- Mahabal, A., Djorgovski, S. G., Donalek, C., Drake, A., Graham, M., Williams, R., Moghaddam, B., & Turmon, M. 2010, in *Astronomical Society of the Pacific Conference Series*, Vol. 434, *Astronomical Data Analysis Software and Systems XIX*, ed. Y. Mizumoto, K.-I. Morita, & M. Ohishi, 115
- Mahabal, A., et al. 2008a, *Astronomische Nachrichten*, 329, 288
- Mahabal, A., et al. 2008b, in *American Institute of Physics Conference Series*, Vol. 1082, *American Institute of Physics Conference Series*, ed. C. A. L. Bailer-Jones, 287–293
- Mahabal, A. A., et al. 2011, *Bulletin of the Astronomical Society of India*, 39, 387
- Manchester, R. N., Hobbs, G. B., Teoh, A., & Hobbs, M. 2005, *aj*, 129, 1993
- Markoff, S., et al. 2005, *ApJ*, 635, 1203
- Matsumura, N., et al. 2007, *AJ*, 133, 1441
- . 2009, *AJ*, 138, 787

- Mauch, T., Murphy, T., Buttery, H. J., Curran, J., Hunstead, R. W., Piestrzynski, B., Robertson, J. G., & Sadler, E. M. 2003, *MNRAS*, 342, 1117
- McClintock, J. E., & Remillard, R. A. 2006, *Cambridge Astrophysics Series*, 39, 157
- McLaughlin, M. A., et al. 2006, *Nature*, 439, 817
- Merloni, A., Heinz, S., & di Matteo, T. 2003, *MNRAS*, 345, 1057
- Metzger, B. D., & Berger, E. 2012, *ApJ*, 746, 48
- Metzger, B. D., Giannios, D., & Mimica, P. 2012, *MNRAS*, 420, 3528
- Miller-Jones, J. C. A., et al. 2012, *MNRAS*, 419, L49
- Mills, B. Y. 1981, *PASA*, 4, 156
- Mirabel, I. F., & Rodríguez, L. F. 1994, *Nature*, 371, 46
- Mitchell, D. A., Robertson, J. G., & Sault, R. J. 2005, *AJ*, 130, 2424
- Moesta, P., Alic, D., Rezzolla, L., Zanotti, O., & Palenzuela, C. 2012, *ApJ*, 749, L32
- Murphy, T., Mauch, T., Green, A., Hunstead, R. W., Piestrzynska, B., Kels, A. P., & Sztajer, P. 2007, *MNRAS*, 382, 382
- Myers, S. T., et al. 2003, *MNRAS*, 341, 1
- Nakar, E., & Piran, T. 2011, arXiv 1102.1020
- Narayan, R. 1996, *ApJ*, 462, 136
- Niinuma, K., et al. 2007, *ApJ*, 657, L37
- . 2009, *ApJ*, 704, 652
- Norris, R. P., et al. 2006, *AJ*, 132, 2409
- . 2011, *PASA*, 28, 215
- Ofek, E. O., Frail, D. A., Breslauer, B., Kulkarni, S. R., Chandra, P., Gal-Yam, A., Kasliwal, M. M., & Gehrels, N. 2011, *ApJ*, 740, 65
- Offringa, A. R., de Bruyn, A. G., & Zaroubi, S. 2012, *MNRAS*, 422, 563
- O'Shaughnessy, R., Kaplan, D. L., Sesana, A., & Kamble, A. 2011, *ApJ*, 743, 136
- Osten, R. A., & Wolk, S. J. 2009, *ApJ*, 691, 1128
- Palenzuela, C., Lehner, L., & Liebling, S. L. 2010, *Science*, 329, 927
- Panagia, N., Van Dyk, S. D., Weiler, K. W., Sramek, R. A., Stockdale, C. J., & Murata, K. P. 2006, *ApJ*, 646, 369
- Pestalozzi, M. R., Benz, A. O., Conway, J. E., & Güdel, M. 2000, *A&A*, 353, 569
- Piran, T., Nakar, E., & Rosswog, S. 2012, arXiv 1204.6242
- Quinlan, J. R. 1986, *Machine Learning*, 1, 81
- Racusin, J. L., et al. 2009, *ApJ*, 698, 43
- Rebbapragada, U., Lo, K. K., Wagstaff, K. L., Reed, C., & Murphy, T. 2012, *VAST Memo* 5
- Rees, M. J. 1990, *Science*, 247, 817
- Rhoads, J. E. 1997, *ApJ*, 487, L1
- Richards, J. W., Starr, D. L., Miller, A. A., Bloom, J. S., Butler, N. R., Brink, H., & Crellin-Quick, A. 2012, arXiv 1204.4180
- Richards, J. W., et al. 2011, *ApJ*, 733, 10
- Rickett, B. J., Kedziora-Chudczer, L., & Jauncey, D. L. 2002, *ApJ*, 581, 103
- Riess, A. G., et al. 2004, *ApJ*, 607, 665
- Rodriguez, C., Taylor, G. B., Zavala, R. T., Peck, A. B., Pollack, L. K., & Romani, R. W. 2006, *ApJ*, 646, 49
- Romani, R., Blandford, R., & Cordes, J. 1987, *Nature*, 328, 324
- Roy, S., Hyman, S. D., Pal, S., Lazio, T. J. W., Ray, P. S., & Kassim, N. E. 2010, *ApJ*, 712, L5
- Ryder, S. D., et al. 2004, *MNRAS*, 349, 1093
- Sari, R., & Piran, T. 1995, *ApJ*, 455, L143
- Sari, R., Piran, T., & Halpern, J. P. 1999, *ApJ*, 519, L17
- Saxton, C. J., Soria, R., Wu, K., & Kuin, N. P. M. 2012, *MNRAS*, 2701
- Schnittman, J. D. 2011, *Classical and Quantum Gravity*, 28, 094021
- Servillat, M., Farrell, S. A., Lin, D., Godet, O., Barret, D., & Webb, N. A. 2011, *ApJ*, 743, 6
- Sesana, A., Gair, J., Berti, E., & Volonteri, M. 2011, *Phys. Rev. D*, 83, 044036
- Shivvers, I., & Berger, E. 2011, *ApJ*, 734, 58
- Siemion, A. P. V., et al. 2012, *ApJ*, 744, 109
- Smith, K. L., Shields, G. A., Bonning, E. W., McMullen, C. C., Rosario, D. J., & Salviander, S. 2010, *ApJ*, 716, 866
- Smolders, B., & Hampson, G. 2002, *IEEE Antennas Propagation Magazine*, 44, 13
- Soderberg, A. M. 2007, in *American Institute of Physics Conference Series*, Vol. 937, *Supernova 1987A: 20 Years After: Supernovae and Gamma-Ray Bursters*, ed. S. Immler, K. Weiler, & R. McCray, 492–499



- Soderberg, A. M., et al. 2010, *Nature*, 463, 513
- Spangler, S., Fanti, R., Gregorini, L., & Padrielli, L. 1989, *A&A*, 209, 315
- Swinbank, J. 2007, in *Bursts, Pulses and Flickering: Wide-Field Monitoring of the Dynamic Radio Sky*
- Thorstensen, J. R., & Armstrong, E. 2005, *AJ*, 130, 759
- Tingay, S. J., et al. 1995, *Nature*, 374, 141
- . 2012, *PASA*, submitted (arXiv 1206.6945)
- Trigilio, C., Leto, P., Umana, G., Buemi, C. S., & Leone, F. 2011, *ApJ*, 739, L10
- Triglio, C., Leto, P., Umana, G., Buemi, C. S., & Leone, F. 2008, *MNRAS*, 384, 1437
- Trott, C. M., Wayth, R. B., Macquart, J.-P. R., & Tingay, S. J. 2011, *ApJ*, 731, 81
- Walker, M. A. 2007, in *Astronomical Society of the Pacific Conference Series*, Vol. 365, *SINS - Small Ionized and Neutral Structures in the Diffuse Interstellar Medium*, ed. M. Haverkorn & W. M. Goss, 299
- Wang, Y.-Y., Wang, L., Xiang, S.-P., Wang, Y., Hao, J.-M., & Yuan, Y.-F. 2010, *Research in Astronomy and Astrophysics*, 10, 199
- Webb, N. A., et al. 2012, *Science*, in press (10.1126/science.1222779)
- Weiler, K. W., Panagia, N., Montes, M. J., & Sramek, R. A. 2002, *ARA&A*, 40, 387
- Welch, J., et al. 2009, *IEEE Proceedings*, 97, 1438
- White, R. L., Becker, R. H., Helfand, D. J., & Gregg, M. D. 1997, *ApJ*, 475, 479
- Whiting, M. T. 2012, *MNRAS*, 421, 3242
- Wiersema, K., et al. 2012, *MNRAS*, 2540
- Yuan, F., et al. 2005, *ApJ*, 620, 905
- Zauderer, B. A., et al. 2011, *Nature*, 476, 425